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## Etude de quelques problèmes aux limites dans des domaines non réguliers

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## ABSTRACT

This thesis is designed to give some results of existence and uniqueness of solutions for some parabolic equations posed in unbounded in time non-cylindrical domains. We give sufficient conditions on the functions of the parametrization of the non regular domains and on the coefficients of the equations under which our problems admit unique solutions. We study the global regularity problem in a suitable parabolic Sobolev space. The method used to prove our main results is based on the technique of the decomposition of domains.

Key words. parabolic equations, heat equation, non-rectangular domains, conical domains, unbounded domains, Dirichlet-Robin conditions, anisotropic Sobolev spaces.

## RÉSUMÉ

Cette thèse a pour but de donner des résultats d'existence et d'unicité de solutions pour certaines équations paraboliques posées dans des domaines non cylindriques et non bornés en temps. Nous donnons des conditions suffisantes sur les fonctions de paramétrisation des domaines non réguliers et sur les coefficients des équations sous lesquelles nos problèmes admettent une solution unique. Nous étudions le problème de régularité globale dans un espace de Sobolev parabolique approprié. La méthode utilisée pour démontrer nos principaux résultats est basée sur la technique de la décomposition de domaines.

Mots clés. Equations paraboliques, équation de la chaleur, domaines non rectangulaires, domaines coniques, domaines non bornés, conditions de Dirichlet-Robin, espaces de Sobolev anisotropes.

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## INTRODUCTION

In this thesis, we shall be concerned by the existence, the uniqueness and maximal regularity of solutions of the following equation :

$$
\begin{equation*}
\partial_{t} u-c(t) \Delta u=f, \tag{1}
\end{equation*}
$$

where $\Delta$ is the Laplacian operator and $c(t)$ is a time dependent coefficient. The second member $f$ belongs to the Lebesgue space of square integrable functions $L^{2}$. The equation (1) is posed in unbounded time-varying domains of $\mathbb{R}^{n}, n=2$ or $n=3$ and it is associated with boundary conditions of the Dirichlet-Robin type or of the Cauchy-Dirichlet type.

The main pecularities of the problems studied in this thesis are the unboundedness of the domains and the fact that they are non regular. Besides being interesting in themselves, such kind of problems are of interest in several fields, see for example [6], [16] and the references therein.

This thesis consists of three chapters. Let us briefly indicate the contents of each chapter.
Chapter 1 is a preliminary chapter in which we recalled essential notions and results that will be used throughout this work. First, we recall some definitions concerning some functional spaces, notably the anisotropic Sobolev spaces. We then prove some technical lemmas. Finally, we present some results on some model parabolic problems that we need to develop further arguments.

In Chapter 2, we will prove well-posedness and regularity results for a one-dimensional parabolic equation, subject to Dirichlet-Robin type boundary conditions and posed in an unbounded in time non-rectangular domain. More precisely, we are concerned by the following problems :

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \in L^{2}(\Omega)  \tag{2}\\
\left.u\right|_{\Gamma_{1}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \in L^{2}(\Omega)  \tag{3}\\
\left.v\right|_{\Gamma_{2}}=0 \\
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=0
\end{array}\right.
$$

where

$$
\Omega:=\left\{(t, x) \in \mathbb{R}^{2}: t>0, \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

with $\varphi_{i} \in C\left(\left[0,+\infty[) \cap C^{1}(0,+\infty), i=1,2\right.\right.$,

$$
\varphi(t):=\varphi_{2}(t)-\varphi_{1}(t)>0 \quad \forall t>0, \text { and } \varphi(0)=0 .
$$

The lateral boundaries of $\Omega$ are defined by

$$
\Gamma_{i}=\left\{\left(t, \varphi_{i}(t)\right) \in \mathbb{R}^{2}: t>0\right\}, i=1,2
$$

The coefficient $c$ is a continuous real-valued function defined on $[0,+\infty[$, differentiable on $] 0,+\infty[$ and such that

$$
0<\alpha \leq c(t) \leq \beta
$$

for every $t \in\left[0,+\infty\left[\right.\right.$, where $\alpha$ and $\beta$ are positive constants. Here, the coefficient $\beta_{i}, i=1,2$ in boundary conditions are real numbers.

Problems (2) and (3) modelize, for instance, the lateral diffusion of two pollutants in a flow of a river with variable width. Note that the Robin type conditions

$$
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
$$

mean for instance, that the flux of diffusion of the pollutants are proportional to their propagations along the wide of the river. The most interesting points of the parabolic problems studied here is the unboundedness of $D$ with respect to the time variable $t$ and the fact that $D$ shrinks at $t=0(\varphi(0)=0)$ which prevent one using the methods in [20] and [21]. It is well known that there are two main approaches for the study of boundary value problems in such non-regular domains. The analysis can be done in weighted spaces with the weight controlling the behavior of the solutions near the singularity of the boundary of the domain (see, for instance, [17], [18] and [19]). Our approach is different. Indeed, the space

$$
\mathcal{H}^{1,2}(\Omega):=\left\{w \in L^{2}(\Omega): \partial_{t} w, \partial_{x} w, \partial_{x}^{2} w \in L^{2}(\Omega)\right\}
$$

used here has low smoothness but one must add assumptions on the type of the domain $\Omega$, as well as conditions on the coefficients $c$ and $\beta_{i}, i=1,2$, near the singular point 0 and in the neighborhood of $+\infty$.

In Chapter 3, we will prove well-posedness and regularity results for a bi-dimensional parabolic equation, subject to Cauchy-Dirichlet boundary conditions and posed in an unbounded in time conical domain of $\mathbb{R}^{3}$. More precisely, we are concerned by the following problem :

$$
\left\{\begin{array}{l}
\partial_{t} w-c(t)\left(\partial_{x}^{2} w+\partial_{y}^{2} w\right)=f \in L^{2}(D)  \tag{4}\\
\left.w\right|_{\partial D}=0
\end{array}\right.
$$

where the conical domain $D$ is defined by

$$
D=\left\{(t, x, y) \in \mathbb{R}^{3}: t>0 ; 0 \leq \sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

where $\varphi \in C\left(\left[0,+\infty[) \cap C^{1}(0,+\infty)\right.\right.$,

$$
\varphi(t)>0 \quad \forall t>0, \text { and } \varphi(0)=0
$$

The coefficient $c$ is a continuous real-valued function defined on $[0,+\infty[$, differentiable on $] 0,+\infty[$ and such that

$$
0<\gamma \leq c(t) \leq \delta
$$

for every $t \in[0,+\infty[$, where $\gamma$ and $\delta$ are positive constants. by using the domain decomposition method, we prove that there exists a unique solution $w$ of Problem (4) belonging to the anisotropic Sobolev space

$$
\mathcal{H}^{1,2}(D)=\left\{w \in H^{1}(D): \partial_{x}^{2} w \in L^{2}(D), \partial_{y}^{2} w \in L^{2}(D), \partial_{x y}^{2} w \in L^{2}(D)\right\}
$$

where $H^{1}(D)$ stands for the Sobolev space defined by

$$
H^{1}(D)=\left\{w \in L^{2}(D): \partial_{t} w \in L^{2}(D), \partial_{x} w \in L^{2}(D), \partial_{y} w \in L^{2}(D)\right\}
$$

Note that results concerning the bounded in time case are obtained for Problem (4) in [13].
We end this thesis by a conclusion and prospects.

## CHAPTER 1

## PRELIMINARIES

The objective of this chapter is to recall the essential notions and results used throughout this work. First, we recall some definitions concerning some functional spaces, notably the anisotropic Sobolev spaces. We then prove some technical lemmas. Finally, we present some results on some model parabolic problems. For more details, see [5], [8] and [21].

### 1.1 Some functional spaces

### 1.1.1 $\quad L^{p}$ spaces

Definition 1.1.1. Let $\Omega$ be an open subset of $\mathbb{R}^{N}$ and $p \in[1,+\infty]$. The space $L^{p}$ is the vectorial space of classes of functions $u$ from $\Omega$ into $\mathbb{R}$, Lebesgue measurables, such that

1. if $1 \leq p<+\infty, \int_{\Omega}|u(x)|^{p} d x<+\infty$,
2. if $p=+\infty, \underset{x \in \Omega}{\operatorname{ess} \sup }|u(x)|<+\infty$, where

$$
\underset{x \in \Omega}{\operatorname{ess} \sup }|u(x)|=\inf \{M| | u(x) \mid \leq M \text { a. e. } x \in \Omega\} .
$$

Proposition 1.1.1. 1. The mapping $\|\cdot\|$ defined from $\mathrm{L}^{p}(\Omega)$ into $\mathbb{R}_{+}$by

$$
u \longmapsto\left\{\begin{array}{lc}
\|u\|_{p}=\left(\int_{\Omega}|u(x)|^{p} d x\right)^{\frac{1}{p}}, & 1 \leq p<+\infty \\
\|u\|_{\infty}=\operatorname{supess}_{x \in \Omega}|u(x)|, & p=\infty,
\end{array}\right.
$$

define a norm on $\mathrm{L}^{p}(\Omega)$, which makes it a Banach space (and a Hilbert space if $p=2$ ).
2. For every real $p \in\left[1,+\infty\left[\right.\right.$, the dual $\mathrm{L}^{p^{\prime}}(\Omega)$ of $\mathrm{L}^{p}(\Omega)$ is isomorph to $\mathrm{L}^{p}(\Omega)$ with $\frac{1}{p}+\frac{1}{p^{\prime}}=1$. The dual mapping is defined by

$$
\mathrm{L}^{p}(\Omega) \times \mathrm{L}^{p^{\prime}}(\Omega) \longmapsto \mathbb{R},(u, v) \longmapsto \int_{\Omega} u(x) v(x) d x,
$$

for each real $p \in\left[1,+\infty\left[\right.\right.$. The bi-dual of $\mathrm{L}^{p}(\Omega)$ can be identified to $\mathrm{L}^{p}(\Omega)$. We said that $\mathrm{L}^{p}(\Omega)$ is a reflexive space.

## Theorem 1.1.1. (Hölder inequality)

Assume that $f \in \mathrm{~L}^{p}$ and $g \in \mathrm{~L}^{p^{\prime}}$ with $\frac{1}{p}+\frac{1}{p^{\prime}}=1$. Then $f g \in \mathrm{~L}^{1}$ and

$$
\int_{\Omega}|f g| d x \leq\|f\|_{L^{p}}\|g\|_{L^{p^{\prime}}} .
$$

We can find the proof of the previous inequality, for example, in [5] (Theorem 4.6 page 50).
Remark 1.1.1. If $p=p^{\prime}=2$, then we obtain the Cauchy-Schwarz inequality

$$
\int_{\Omega}|f g| d x \leq\|f\|_{\mathrm{L}^{2}}\|g\|_{\mathrm{L}^{2}} .
$$

### 1.1.2 Sobolev spaces

Assume that $\Omega$ is an open domain in $\mathbb{R}^{N}$. For $m \in \mathbb{N}$ and $p \in[1,+\infty[$, the Sobolev space $W^{m, p}(\Omega)$ is defined by :

$$
W^{m, p}(\Omega)=\left\{u: \Omega \longrightarrow \mathbb{R}\left|D^{\alpha} u \in L^{p}(\Omega), \forall \alpha \in \mathbb{N}^{N}, 0 \leq|\alpha| \leq m\right\}\right.
$$

where for any $\alpha=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{N}\right) \in \mathbb{N}^{N}$ we note $|\alpha|=\alpha_{1}+\alpha_{2}+\ldots+\alpha_{N}$ and $D^{\alpha} u=$ $\partial_{x_{1}}^{\alpha_{1}} \partial_{x_{2}}^{\alpha_{2}} \ldots \partial_{x_{N}}^{\alpha_{N}} u$. The space $W^{m, p}(\Omega)$ is equipped with the norm

$$
\begin{equation*}
\|u\|_{W^{m, p}(\Omega)}=\sum_{0 \leq|\alpha| \leq m}\left\|D^{\alpha} u\right\|_{L^{p}(\Omega)}^{p} \tag{1.1}
\end{equation*}
$$

or with the equivalent norm

$$
\|u\|_{W^{m, p}(\Omega)}=\left(\sum_{0 \leq|\alpha| \leq m}\left\|D^{\alpha} u\right\|_{L^{p}(\Omega)}^{p}\right)^{\frac{1}{p}} .
$$

Proposition 1.1.2. $W^{m, p}(\Omega)$ is a separable Banach space. It is reflexive for $1<p<\infty$.
Remark 1.1.2. 1. If $p=2$, we usually write $W^{m, 2}(\Omega)=\mathcal{H}^{m}(\Omega)$. When equipped with the inner product

$$
(u \mid v)_{\mathcal{H}^{m}(\Omega)}=\sum_{0 \leq|\alpha| \leq m}\left(D^{\alpha} u \mid D^{\alpha} v\right)_{\mathrm{L}^{2}(\Omega)},
$$

$\mathcal{H}^{m}(\Omega)$ is a Hilbert space.
2. If $\Omega$ is an open bounded set with smooth boundary $\Gamma$, the the norm (1.1) is equivalente to the following norm

$$
\|u\|_{L^{p}(\Omega)}+\sum_{|\alpha| \leq m}\left\|D^{\alpha} u\right\|_{L^{p}(\Omega)} .
$$

Definition 1.1.2. $D(\Omega)$ is the set of functions of class $C^{\infty}(\Omega)$ with compact support in $\Omega$, $i$. $e$.

$$
\mathrm{D}(\Omega)=\left\{u: \Omega \rightarrow \mathbb{R} ; u \in C^{\infty}(\Omega) \text { et } \operatorname{supp}(u) \subset K \subset \Omega, K \text { compact }\right\} .
$$

Definition 1.1.3. We note by $W_{0}^{m, p}(\Omega)=\overline{\mathrm{D}(\Omega)}{ }^{W^{m, p}(\Omega)}$ the closure of de $\mathrm{D}(\Omega)$ in $W^{m, p}(\Omega)$.
Theorem 1.1.2. $D(\bar{\Omega})$ is dense in $W^{m, p}(\Omega)$.

In the sequel, $W^{1, p}(\Omega)$ will be equipped with the norm

$$
\|u\|_{W^{1, p}(\Omega)}=\left(\|u\|_{\mathrm{L}^{P}(\Omega)}^{p}+\sum_{j=1}^{N}\left\|\frac{\partial u}{\partial x_{j}}\right\|_{\mathrm{L}^{P}(\Omega)}^{p}\right)^{\frac{1}{p}}
$$

or with the equivalent norm

$$
\|u\|_{W^{1, p}(\Omega)}=\|u\|_{L^{p}(\Omega)}+\sum_{j=1}^{N}\left\|\frac{\partial u}{\partial x_{j}}\right\|_{L^{p}(\Omega)} .
$$

In the case of a bounded set $\Omega, W^{1, p}(\Omega)$ can be equipped with the following norm:

$$
\|u\|_{W_{0}^{1, p}(\Omega)}=\left(\int_{\Omega}|\nabla u|^{p} d x\right)^{\frac{1}{p}}
$$

which is called the gradient norm. The equivalence of these norms can be obtained by using the following result :

## Theorem 1.1.3. (Poincaré Inequality)

Assume that $\Omega$ is bounded. Then, there exists a constant $C$ (which depends on $\Omega$ and on $p$ ) such that

$$
\|u\|_{\mathrm{L}^{p}(\Omega)} \leq C\|\nabla u\|_{\mathrm{L}^{p}(\Omega)}, \forall u \in W_{0}^{1, p}(\Omega), 1 \leq p<\infty .
$$

The mapping $u \longmapsto\|\nabla u\|_{L^{p}(\Omega)}$ is a norm on $W_{0}^{1, p}(\Omega)$ which is equivalente to the norm induced by $\|\cdot\|_{W^{1, p}(\Omega)}$.

## Theorem 1.1.4. (Green Formula)

Assume that $\Omega$ is a open bounded set of $\mathbb{R}^{N}$ of class $C^{1}$ such that its boundary is bounded. Then, for each $u \in \mathcal{H}^{2}(\Omega)$ and for each $v \in \mathcal{H}^{1}(\Omega)$, we have

$$
\int_{\Omega} \Delta u \cdot v d x=-\int_{\Omega} \nabla u \cdot \nabla v d x+\int_{\Gamma} \frac{\partial u}{\partial n} \cdot v d \sigma
$$

where $\frac{\partial u}{\partial n}=\nabla u \cdot n=\sum_{i=1}^{N} \frac{\partial u}{\partial x_{i}} \eta_{i}$ is the normal derivative and $n={ }^{t}\left(n_{1}, \ldots, n_{N}\right)$ is the unit normal vector.

The Green formula can be given by the following theorem :

## Theorem 1.1.5. (Green Formula)

Assume that $\Omega$ is an open bounded domain of $\mathbb{R}^{N}$ with Lipschitz boundary $\Gamma=\partial \Omega$ and $1<p<$ $\infty$. Then, for every $u \in W^{1, p}(\Omega)$ and for every $v \in W^{-1, p^{\prime}}(\Omega)$, we have

$$
\begin{equation*}
\int_{\Omega} \frac{\partial u}{\partial x_{i}} v d x=-\int_{\Omega} u \frac{\partial v}{\partial x_{i}} d x+\int_{\Gamma} u v v_{i} d \sigma, i=1 \ldots N, \tag{1.2}
\end{equation*}
$$

where $\frac{\partial u}{\partial x_{i}}$ and $\frac{\partial v}{\partial x_{i}}$ are taken in the distributional sens, $v_{i}$ is the $i$-th composante of unit vector of the normale to $\partial \Omega$. Here, $W^{-1, p^{\prime}}(\Omega)$ denotes the dual space of $W^{1, p}(\Omega), 1<p<\infty$.

### 1.1.3 Anisotropic Sobolev spaces

It is well known that Lebesgue and Sobolev spaces are the essential tools of functional analysis for the study of partial differential equations. The study of the regularity of solutions can be carried out by their belonging to some functional spaces. In this section, we introduce the so-called anisotropic Sobolev spaces $\mathcal{H}^{1,2}$ built on the Lebesgue space of square integrable
functions $L^{2}$. These spaces are the natural ones adopted in the study of second-order parabolic equations. The main features of the spaces $\mathcal{H}^{1,2}$ is that smoothness with respect to spatial variables are twice as high with respect to time. For more details and proofs we refer to Lions and Magenes [21].

Let us introduce some notations. Let $\Omega$ be a bounded open subset of $\mathbb{R}^{n}$. We denote $Q=(0, T) \times \Omega$ with $T$ maybe finite or infinite. $L^{2}(Q)$ denotes the space of (class of) functions $u$ which are square integrable on $Q$ with the Lebesgue measure $d t d x$, such that

$$
\|u\|_{L^{2}(Q)}=\left(\int_{Q}|u|^{2} d t d x\right)^{1 / 2}<\infty
$$

Now we define the anisotropic Sobolev space $\mathcal{H}^{1,2}(Q)$
Definition 1.1.4. $\mathcal{H}^{1,2}(Q)=\left\{u \in L^{2}(Q), \partial_{t} u \in L^{2}(Q), \partial^{\alpha} u \in L^{2}(Q)\right.$, for each $\left.|\alpha| \leq 2\right\}$.

$$
\partial^{\alpha} u=\partial_{x_{1}}^{\alpha_{1}} \ldots \partial_{x_{n}}^{\alpha_{n}} u, \alpha=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{n}\right),|\alpha|=\alpha_{1}+\ldots+\alpha_{n} \leq 2 . \text { Of course, it is a Hilbert }
$$ space with inner product

$$
(u, v)_{\mathcal{H}^{1,2}(Q)}=\left(\partial_{t} u, \partial_{t} v\right)_{L^{2}(Q)}+\sum_{|\alpha| \leq 2}\left(\partial^{\alpha} u, \partial^{\alpha} v\right)_{L^{2}(Q)}
$$

$\mathcal{H}^{1,2}(Q)$ is equipped with the norm

$$
\|u\|_{\mathcal{H}^{1,2}(Q)}=\left(\left\|\partial_{t} u\right\|_{L^{2}(Q)}^{2}+\sum_{|\alpha| \leq 2}\left\|\partial^{\alpha} u\right\|_{L^{2}(Q)}^{2}\right)^{1 / 2}
$$

Remark 1.1.3. The spaces $H^{1,2}(Q)$, are said to be anisotropic in the sense that orders of differentiability in the time and spatial directions are not equal.

### 1.2 Technical Lemmas

The following result is well known, see for example [21].
Lemma 1.2.1. Let $D(0,1)$ be the unit disk of $\mathbb{R}^{2}$. Then, the Laplace operator $\Delta=\partial_{x}^{2}+\partial_{y}^{2}$ : $H^{2}(D(0,1)) \cap H_{0}^{1}(D(0,1)) \longrightarrow L^{2}(D(0,1))$ is an isomorphism. Moreover, there exists a constant $C>0$ such that

$$
\|v\|_{H^{2}(D(0,1))} \leq C\|\Delta v\|_{L^{2}(D(0,1))}, \forall v \in H^{2}(D(0,1)) \cap H_{0}^{1}(D(0,1)) .
$$

Here, $H^{2}$ and $H_{0}^{1}$ are the usual Sobolev spaces defined, for instance, in Lions-Magenes [21].

In the sequel, assume that $\beta_{i}, i=1,2$ are real numbers such that

$$
\begin{equation*}
(-1)^{i} \beta_{i}>0, i=1,2 . \tag{1.3}
\end{equation*}
$$

Lemma 1.2.2. Assume that $\beta_{i}, i=1,2$ fulfil the condition (1.3). Then, there exists a positive constant $K_{1}$ such that for each $(u, v) \in H_{\gamma}^{2}(0,1) \times H_{\delta}^{2}(0,1)$

$$
\begin{aligned}
& \left\|u^{(k)}\right\|_{L^{2}(0,1)} \leq K_{1}\left\|u^{(2)}\right\|_{L^{2}(0,1)}, k=0,1, \\
& \left\|v^{(k)}\right\|_{L^{2}(0,1)} \leq K_{1}\left\|v^{(2)}\right\|_{L^{2}(0,1)}, k=0,1,
\end{aligned}
$$

where

$$
H_{\gamma}^{2}(0,1)=\left\{u \in H^{2}(0,1): u^{\prime}(1)+\beta_{2} u(1)=u(0)=0\right\}
$$

and

$$
H_{\delta}^{2}(0,1)=\left\{v \in H^{2}(0,1): v^{\prime}(0)+\beta_{1} v(0)=v(1)=0\right\} .
$$

Proof. Let $h_{1}, h_{2}$ be arbitrary fixed elements of $L^{2}(0,1)$. Every solution of the ordinary differential equation $u^{\prime \prime}=h_{1}$, (respectively, $v^{\prime \prime}=h_{2}$,) is of the form

$$
u(y)=\int_{0}^{y}\left\{\int_{0}^{x} h_{1}(s) d s\right\} d x+y u^{\prime}(0)+u(0), y \in[0,1]
$$

(respectively,

$$
\left.v(y)=\int_{0}^{y}\left\{\int_{0}^{x} h_{2}(s) d s\right\} d x+y v^{\prime}(0)+v(0), y \in[0,1]\right)
$$

The variables $u(0)$ and $u^{\prime}(0)$ (respectively, $v(0)$ and $v^{\prime}(0)$ ) are to be determined in a unique way such that the boundary conditions $u^{\prime}(1)+\beta_{2} u(1)=u(0)=0$ (respectively, $v^{\prime}(0)+\beta_{1} v(0)=$ $v(1)=0)$ are satisfied.

From the preceding representation of the solution (and thus also its derivative) and from the required boundary conditions we obtain the following system to be solved:

$$
\left\{\begin{array}{l}
\left(1+\beta_{2}\right) u^{\prime}(0)+\beta_{2} u(0)=-\int_{0}^{1} h_{1}(s) d s-\beta_{2} \int_{0}^{1}\left\{\int_{0}^{x} h_{1}(s) d s\right\} d x \\
0 u^{\prime}(0)+u(0)=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
v^{\prime}(0)+v(0)=-\int_{0}^{1}\left\{\int_{0}^{x} h_{2}(s) d s\right\} d x \\
\left.v^{\prime}(0)+\beta_{1} v(0)=0\right)
\end{array}\right.
$$

This system in the unknowns $u(0)$ and $u^{\prime}(0)$ (respectively, $v(0)$ and $\left.v^{\prime}(0)\right)$ is uniquely solvable if and only if

$$
\beta_{2}+1 \neq 0,
$$

(respectively,

$$
\left.\beta_{1}-1 \neq 0\right) .
$$

This condition is verified thanks to (1.3). Finally, the unique solution of the problem

$$
\left\{\begin{array}{l}
u^{\prime \prime}=h_{1} \\
u(0)=0 \\
u^{\prime}(1)+\beta_{2} u(1)=0,
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
v^{\prime \prime}=h_{2} \\
v(1)=0 \\
\left.v^{\prime}(0)+\beta_{1} v(0)=0\right)
\end{array}\right.
$$

is given by

$$
u(y)=\int_{0}^{y}\left\{\int_{0}^{x} h_{1}(s) d s\right\} d x+y u^{\prime}(0)
$$

(respectively,

$$
\left.v(y)=\int_{0}^{y}\left\{\int_{0}^{x} h_{2}(s) d s\right\} d x+y v^{\prime}(0)+v(0)\right)
$$

where

$$
u^{\prime}(0)=\frac{-\int_{0}^{1} h_{1}(s) d s-\beta_{2} \int_{0}^{1}\left\{\int_{0}^{x} h_{1}(s) d s\right\} d x}{\beta_{2}+1}
$$

(respectively,

$$
\left\{\begin{array}{l}
v(0)=\frac{-\int_{0}^{1}\left\{\int_{0}^{x} h_{2}(s) d s\right\} d x}{\beta_{1}-1} \\
\left.v^{\prime}(0)=-\beta_{1} v(0)\right) .
\end{array}\right.
$$

Using the Cauchy-Schwarz inequality, we obtain the following estimates

$$
\left|u^{\prime}(0)\right| \leq C\left\|h_{1}\right\|_{L^{2}(0,1)},
$$

(respectively,

$$
\begin{aligned}
& |v(0)| \leq C\left\|h_{2}\right\|_{L^{2}(0,1)} \\
& \left.\left|v^{\prime}(0)\right| \leq C\left\|h_{2}\right\|_{L^{2}(0,1)}\right),
\end{aligned}
$$

which will allow us to obtain the desired estimates.

Lemma 1.2.3. Under the assumption (1.3) on $\beta_{i}, i=1,2$, there exists a positive constant $C_{1}$ (independent of $a$ and b) such that for each $(u, v) \in H_{\gamma}^{2}(a, b) \times H_{\delta}^{2}(a, b)$

$$
\begin{aligned}
&\left\|u^{(k)}\right\|_{L^{2}(a, b)}^{2} \leq C_{1}(b-a)^{2(2-k)}\left\|u^{(2)}\right\|_{L^{2}(a, b)}^{2}, k=0,1, \\
&\left\|v^{(k)}\right\|_{L^{2}(a, b)}^{2} \leq C_{1}(b-a)^{2(2-k)}\left\|v^{(2)}\right\|_{L^{2}(a, b)}^{2}, k=0,1,
\end{aligned}
$$

where,

$$
\begin{aligned}
& H_{\gamma}^{2}(a, b)=\left\{u \in H^{2}(a, b): u(a)=0, u^{\prime}(b)+\frac{\beta_{2}}{b-a} u(b)=0\right\}, \\
& H_{\delta}^{2}(a, b)=\left\{v \in H^{2}(a, b): v^{\prime}(a)+\frac{\beta_{1}}{b-a} v(a)=0, v(b)=0\right\} .
\end{aligned}
$$

Proof. It is a direct consequence of Lemma 1.2.2 by using the following affine change of variable

$$
[0,1] \longrightarrow[a, b], x \mapsto(1-x) a+x b=y .
$$

### 1.3 Some model parabolic problems

The following results are consequences of Theorem 4.3 ([21], Vol.2).
Proposition 1.3.1. Let $R$ be the cylinder $] 0, T\left[\times B(0,1)\right.$ where $B(0,1)$ is the unit disk of $\mathbb{R}^{2}$, $f \in L^{2}(R)$ and $u_{0} \in H^{1}\left(\gamma_{0}\right)$. Then the problem

$$
\left\{\begin{array}{l}
\partial_{t} u-\Delta u=f \text { in } R \\
\left.u\right|_{\gamma_{0}}=u_{0} \\
\left.u\right|_{\gamma_{1}}=0,
\end{array}\right.
$$

where $\Delta=\sum_{j=1}^{2} \partial_{x_{j}}^{2}, \gamma_{0}=\{0\} \times B(0,1)$ et $\left.\gamma_{1}=\right] 0, T[\times \partial B(0,1)$, admits a (unique) solution $u \in H^{1,2}(R)$ if and only if the following compatibility condition is satisfied

$$
\left.u_{0}\right|_{\partial \gamma_{0}}=0 .
$$

Proposition 1.3.2. (Theorem 4.3, [21]) Let $Q=] 0, T[\times] 0,1\left[, f \in L^{2}(Q)\right.$ and $\phi \in H^{1}\left(\gamma_{0}\right)$. Then the following initial/boundary value problems :

$$
\left\{\begin{array}{l}
\partial_{t} u-\partial_{x}^{2} u=f_{1} \in L^{2}(Q) \\
\left.u\right|_{\Gamma_{0}}=\phi \\
\left.u\right|_{\Gamma_{1}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\partial_{t} v-\partial_{x}^{2} v=f_{2} \in L^{2}(Q) \\
\left.v\right|_{\Gamma_{0}}=\phi \\
\left.v\right|_{\Gamma_{2}}=0 \\
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=0
\end{array}\right.
$$

admit unique solutions in $\mathcal{H}^{1,2}(Q)$. Here, $\left.\gamma_{0}=\{0\} \times\right] 0,1\left[, \gamma_{1}=\right] 0, T\left[\times\{0\}\right.$ and $\left.\gamma_{2}=\right] 0, T[\times\{1\}$, The coefficients $\beta_{i}, 1=1,2$ are real numbers satisfying $(-1)^{i} \beta_{i}>0$.

## CHAPTER 2

## GLOBAL IN TIME RESULTS FOR A <br> PARABOLIC EQUATION SOLUTION IN NON-RECTANGULAR DOMAINS

This chapter deals with the parabolic equation

$$
\partial_{t} w-c(t) \partial_{x}^{2} w=f \text { in } D, D=\left\{(t, x) \in \mathbb{R}^{2}: t>0, \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

with $\varphi_{i}:[0,+\infty[\rightarrow \mathbb{R}, i=1,2$ and $c:[0,+\infty[\rightarrow \mathbb{R}$ satisfying some conditions and the problem is supplemented with boundary conditions of Dirichlet-Robin type. We study the global regularity problem in a suitable parabolic Sobolev space. We prove in particular that for $f \in L^{2}(D)$ there exists a unique solution $w$ such that $w, \partial_{t} w, \partial^{j} w \in L^{2}(D), j=1,2$. Notice that the case of bounded non-rectangular domains is studied in [15]. The proof is based on energy estimates after transforming the problem in a strip region combined with some interpolation inequality. This work complements the results obtained in [26] in the case of Cauchy-Dirichlet boundary conditions.

### 2.1 Introduction and statement of the main result

Let $D$ be an open set of $\mathbb{R}^{2}$ defined by

$$
D:=\left\{(t, x) \in \mathbb{R}^{2}: t>0, \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

where $\varphi_{i} \in C\left(\left[0,+\infty[) \cap C^{1}(0,+\infty), i=1,2\right.\right.$,

$$
\varphi(t):=\varphi_{2}(t)-\varphi_{1}(t)>0 \quad \forall t>0, \text { and } \varphi(0)=0 .
$$

The lateral boundaries of $D$ are defined by

$$
\Gamma_{i}=\left\{\left(t, \varphi_{i}(t)\right) \in \mathbb{R}^{2}: t>0\right\}, i=1,2
$$

Let us introduce the following functional space:

$$
\mathcal{H}^{1,2}(D):=\left\{w \in L^{2}(D): \partial_{t} w, \partial_{x} w, \partial_{x}^{2} w \in L^{2}(D)\right\}
$$

where $L^{2}(D)$ stands for the usual Lebesgue space of square-integrable functions on $D$. The space $\mathcal{H}^{1,2}(D)$ is equipped with the natural norm, that is

$$
\|w\|_{\mathcal{H}^{1,2}(D)}^{2}=\|w\|_{L^{2}(D)}^{2}+\left\|\partial_{t} w\right\|_{L^{2}(D)}^{2}+\sum_{j=1}^{2}\left\|\partial_{x}^{j} w\right\|_{L^{2}(D)}^{2}
$$

We consider the problems: to find a function $u \in \mathcal{H}^{1,2}(D)$ (respectively, $v \in \mathcal{H}^{1,2}(D)$ ) that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \text { a.e. on } D \tag{2.1}
\end{equation*}
$$

(respectively,

$$
\begin{equation*}
\left.\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \text { a.e. on } D\right) \tag{2.2}
\end{equation*}
$$

and the boundary conditions

$$
\begin{equation*}
\left.u\right|_{\Gamma_{1}}=\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0, \tag{2.3}
\end{equation*}
$$

(respectively,

$$
\begin{equation*}
\left.\left.v\right|_{\Gamma_{2}}=\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=0\right), \tag{2.4}
\end{equation*}
$$

where $f_{i} \in L^{2}(D), i=1,2$ and the coefficient $c$ is a continuous real-valued function defined on $[0,+\infty[$, differentiable on $] 0,+\infty[$ and such that

$$
0<\alpha \leq c(t) \leq \beta
$$

for every $t \in\left[0,+\infty\left[\right.\right.$, where $\alpha$ and $\beta$ are positive constants. Here, the coefficient $\beta_{i}, i=1,2$ in boundary conditions are real numbers such that


Fig.1: The unbounded non-rectangular domain D.

Problems (2.1)-(2.3) and (2.2)-(2.4) modelize, for instance, the lateral diffusion of two pollutants in a flow of a river with variable width. Note that the Robin type conditions

$$
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
$$

mean for instance, that the flux of diffusion of the pollutants are proportional to their propagations along the wide of the river. The most interesting points of the parabolic problems studied here is the unboundedness of $D$ with respect to the time variable $t$ and the fact that $D$ shrinks at $t=0(\varphi(0)=0)$ which prevent one using the methods in [20] and [21]. It is well known that there are two main approaches for the study of boundary value problems in such non-regular
domains. The analysis can be done in weighted spaces with the weight controlling the behavior of the solutions near the singularity of the boundary of the domain (see, for instance, [17], [18] and [19]). Our approach is different. Indeed, the space $\mathcal{H}^{1,2}$ used here has low smoothness but one must add assumptions on the type of the domain $D$, as well as conditions on the coefficients $c$ and $\beta_{i}, i=1,2$, near the singular point 0 and in the neighborhood of $+\infty$. So, our main result is the following:

Theorem 2.1.1. Let us assume that

$$
\begin{gather*}
\left|\varphi_{i}^{\prime}(t)\right| \varphi(t) \rightarrow 0 \quad \text { as } t \rightarrow 0^{+}, \quad i=1,2,  \tag{2.5}\\
\left.(-1)^{i}\left(2 c(t) \beta_{i}-\varphi_{i}^{\prime}(t)\right) \geq 0 \text { a.e. } t \in\right] 0,+\infty[, i=1,2,  \tag{2.6}\\
\varphi \text { and } \varphi^{\prime} \text { are uniformly bounded in a neighborhood of }+\infty,  \tag{2.7}\\
c \text { is a decreasing function in }] 0,+\infty[, \tag{2.8}
\end{gather*}
$$

and one of the following conditions is satisfied
(a) $\varphi$ is increasing in a neighborhood of $+\infty$,
(b) $\exists M>0:\left|\varphi^{\prime}\right| \varphi \leq M c(t)$.

Then Problem (2.1),(2.3) (respectively, Problem (2.2),(2.4)) admits a unique solution $u \in$ $\mathcal{H}^{1,2}(D)$ (respectively, $v \in \mathcal{H}^{1,2}(D)$ ).

The case where $D$ is bounded (with $c(t)=1$ ) is studied in [15]. The case where $\beta_{1}=\infty$ (or $\left.\beta_{2}=\infty\right)$ corresponding to Cauchy-Dirichlet boundary conditions is studied in [26]. Whereas second-order parabolic equations in bounded non-cylindrical domains are well studied (see for instance [2], [7], [11], [22], [23], [25] and the references therein), the literature concerning unbounded non-cylindrical domains does not seem to be very rich. The regularity of the heat equation solution in a non-smooth and unbounded domain (in the $x$ direction) is obtained in [14], [10] and [3].

In the next sections, we prove Theorem 2.1.1 in four steps:
(1) case of a bounded domain which can be transformed into a rectangle;
(2) case of an unbounded domain which can be transformed into a half strip;
(3) case of a small in time bounded triangular domain;
(4) finally, we use the previous steps and a trace result to complete the proof of Theorem 2.1.1.

### 2.2 The case of a bounded domain which can be transformed into a rectangle

Let $T$ be an arbitrary positive number. Denote by

$$
D_{1}:=\left\{(t, x) \in \mathbb{R}^{2}: 0<t<T ; \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

with $\varphi(t)>0$ for all $t \in[0, T]$ and consider the following problems:

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \text { a.e. on } D_{1},  \tag{2.9}\\
\left.u\right|_{\Gamma_{1}}=\left.u\right|_{\Gamma_{0}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \text { a.e. on } D_{1}  \tag{2.10}\\
\left.v\right|_{\Gamma_{2}}=\left.v\right|_{\Gamma_{0}}=0 \\
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=
\end{array}\right.
$$

where $f_{i} \in L^{2}\left(D_{1}\right), i=1,2$ and $\Gamma_{0}$ is the part of $\partial D_{1}$ where $t=0$.


Fig.2: The bounded domain $D_{1}$.

Let us denote the inner product in $L^{2}\left(D_{1}\right)$ by $\langle.,$.$\rangle . Then, the uniqueness of the solutions may$ be obtained by developing the inner products

$$
\left\langle\partial_{t} u-c(t) \partial_{x}^{2} u, u\right\rangle \text { and }\left\langle\partial_{t} v-c(t) \partial_{x}^{2} v, v\right\rangle .
$$

Indeed, Let us consider $u \in \mathcal{H}^{1,2}\left(D_{1}\right)$ (respectively, $\left.v \in \mathcal{H}^{1,2}\left(D_{1}\right)\right)$ a solution of Problem (2.9) (respectively, of Problem (2.10)) with a null right-hand side terms. So,

$$
\partial_{t} u-c(t) \partial_{x}^{2} u=\partial_{t} v-c(t) \partial_{x}^{2} v=0 \text { in } D_{1} .
$$

In addition $u$ and $v$ fulfil the boundary conditions

$$
\left.u\right|_{\Gamma_{0}}=\left.v\right|_{\Gamma_{0}}=\left.u\right|_{\Gamma_{1}}=\left.v\right|_{\Gamma_{2}}=\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0 .
$$

Using Green formula, we have

$$
\begin{aligned}
& \int_{D_{1}}\left(\partial_{t} u-c(t) \partial_{x}^{2} u\right) u d t d x+\int_{D_{1}}\left(\partial_{t} v-c(t) \partial_{x}^{2} v\right) v d t d x \\
& =\int_{\partial D_{1}}\left(\frac{1}{2}|u|^{2} \nu_{t}-c(t) \partial_{x} u \cdot u \nu_{x}\right) d \sigma+\int_{\partial D_{1}}\left(\frac{1}{2}|v|^{2} \nu_{t}-c(t) \partial_{x} v \cdot v \nu_{x}\right) d \sigma \\
& +\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}\right) d t d x+\int_{D_{1}} c(t)\left(\left|\partial_{x} v\right|^{2}\right) d t d x,
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the unit outward normal vector at $\partial D_{1}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $D_{1}$ where $t=0$, we have $u=v=0$. Accordingly the corresponding boundary integrals vanish. On the part of the boundary of $D_{1}$ where $t=T$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
\frac{1}{2} \int_{\varphi_{1}(T)}^{\varphi_{2}(T)}\left[|u|^{2}(T, x)+|v|^{2}(T, x)\right] d x
$$

is nonnegative. On the parts of the boundary where $x=\varphi_{i}(t), i=1,2$, we have

$$
\nu_{x}=\frac{(-1)^{i}}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}, \nu_{t}=\frac{(-1)^{i+1} \varphi_{i}^{\prime}(t)}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}
$$

and

$$
u\left(t, \varphi_{1}(t)\right)=v\left(t, \varphi_{2}(t)\right)=\partial_{x} u\left(t, \varphi_{2}(t)\right)+\beta_{2} u\left(t, \varphi_{2}(t)\right)=\partial_{x} v\left(t, \varphi_{1}(t)\right)+\beta_{1} v\left(t, \varphi_{1}(t)\right)=0
$$

Consequently, the corresponding integral is

$$
\int_{0}^{T}\left(2 c(t) \beta_{2}-\varphi_{2}^{\prime}(t)\right) u^{2}\left(t, \varphi_{2}(t)\right) d t+\int_{0}^{T}\left(-2 c(t) \beta_{1}+\varphi_{1}^{\prime}(t)\right) v^{2}\left(t, \varphi_{1}(t)\right) d t
$$

Then, we obtain

$$
\begin{aligned}
& \int_{D_{1}}\left(\partial_{t} u-c(t) \partial_{x}^{2} u\right) u d t d x+\int_{D_{1}}\left(\partial_{t} v-c(t) \partial_{x}^{2} v\right) v d t d x \\
& =\int_{0}^{T}\left(2 c(t) \beta_{2}-\varphi_{2}^{\prime}(t)\right) u^{2}\left(t, \varphi_{2}(t)\right) d t+\int_{0}^{T}\left(-2 c(t) \beta_{1}+\varphi_{1}^{\prime}(t)\right) v^{2}\left(t, \varphi_{1}(t)\right) d t \\
& \quad+\frac{1}{2} \int_{\varphi_{1}(T)}^{\varphi_{2}(T)}\left[|u|^{2}(T, x)+|v|^{2}(T, x)\right] d x+\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}\right) d t d x+\int_{D_{1}} c(t)\left(\left|\partial_{x} v\right|^{2}\right) d t d x .
\end{aligned}
$$

Consequently using the fact that $u$ and $v$ are the solutions yields

$$
\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}\right) d t d x+\int_{D_{1}} c(t)\left(\left|\partial_{x} v\right|^{2}\right) d t d x=0
$$

because thanks to the condition (2.6) and to the fact that $c(t)>0$ for every $t \in[0,+\infty[$, we have

$$
\begin{aligned}
& \int_{0}^{T}\left(2 c(t) \beta_{2}-\varphi_{2}^{\prime}(t)\right) u^{2}\left(t, \varphi_{2}(t)\right) d t+\int_{0}^{T}\left(-2 c(t) \beta_{1}+\varphi_{1}^{\prime}(t)\right) v^{2}\left(t, \varphi_{1}(t)\right) d t \\
& +\frac{1}{2} \int_{\varphi_{1}(T)}^{\varphi_{2}(T)}\left[|u|^{2}(T, x)+|v|^{2}(T, x)\right] d x+\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}\right) d t d x \geq 0
\end{aligned}
$$

This implies that $\left|\partial_{x} u\right|^{2}+\left|\partial_{x} v\right|^{2}=0$ and consequently $\partial_{x}^{2} u=\partial_{x}^{2} v=0$. Then, the hypothesis $\partial_{t} u-c(t) \partial_{x}^{2} u=\partial_{t} v-c(t) \partial_{x}^{2} v=0$ gives $\partial_{t} u=\partial_{t} v=0$. Thus, $u, v$ are constants. The boundary conditions and the fact that $\beta_{i} \neq 0, i=1,2$ imply that $u=v=0$ in $D_{1}$. This proves the uniqueness of the solutions of Problems (2.9) and (2.10).

Now, let us look at the existence of solutions for Problems (2.9) and (2.10). The change of variables $(t, x)$ to $\left(t, \frac{x-\varphi_{1}(t)}{\varphi(t)}\right)$ transforms $D_{1}$ into the rectangle $\left.Q=\right] 0, T[\times] 0,1[$ and Problem (2.9) (respectively, Problem (2.10)) becomes the following:

$$
\left\{\begin{array}{l}
\partial_{t} u+a(t, x) \partial_{x} u-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u=f_{1} \text { a.e. on } Q \\
\left.u\right|_{t=0}=\left.u\right|_{x=0}=0 \\
\partial_{x} u+\left.\beta_{2} \varphi(t) u\right|_{x=1}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v+a(t, x) \partial_{x} v-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v=f_{2} \text { a.e. on } Q \\
\left.v\right|_{t=0}=\left.v\right|_{x=1}=0 \\
\left.\partial_{x} v+\left.\beta_{1} \varphi(t) v\right|_{x=0}=0\right)
\end{array}\right.
$$

where $f_{i} \in L^{2}(Q), i=1,2$ and $a(t, x)=-\frac{x \varphi^{\prime}(t)+\varphi_{1}^{\prime}(t)}{\varphi(t)}$. Observe that the coefficient $a$ is bounded. So, the operator

$$
a(t, x) \partial_{x}: \mathcal{H}^{1,2}(Q) \longrightarrow L^{2}(Q)
$$

is compact. Hence, it is sufficient to study the following problem:

$$
\left\{\begin{array}{l}
\partial_{t} u-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u=f_{1} \text { a.e. on } Q  \tag{2.11}\\
\left.u\right|_{t=0}=\left.u\right|_{x=0}=0 \\
\partial_{x} u+\left.\beta_{2} \varphi(t) u\right|_{x=1}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v=f_{2} \text { a.e. on } Q  \tag{2.12}\\
\left.v\right|_{t=0}=\left.v\right|_{x=1}=0 \\
\left.\partial_{x} v+\left.\beta_{1} \varphi(t) v\right|_{x=0}=0\right)
\end{array}\right.
$$

where $f_{i} \in L^{2}(Q), i=1,2$. It is clear that Problem (2.11) (respectively, Problem (2.12)) admits a (unique) solution $u \in \mathcal{H}^{1,2}(Q)$ (respectively, $v \in \mathcal{H}^{1,2}(Q)$ ) because the coefficient $\frac{c(t)}{\varphi^{2}(t)}$ satisfies the "uniform parabolicity" condition (see, for example [1]). On other hand, it is easy to verify that the aforementioned change of variable conserves the spaces $L^{2}$ and $\mathcal{H}^{1,2}$. Consequently, we have proved the following theorem:

Theorem 2.2.1. Problem (2.9) (respectively, Problem (2.10)) admits a (unique) solution $u \in$ $\mathcal{H}^{1,2}\left(D_{1}\right)$ (respectively, $v \in \mathcal{H}^{1,2}\left(D_{1}\right)$ ).

### 2.3 The case of an unbounded domain which can be transformed into a half strip

In this case, we set

$$
D_{2}:=\left\{(t, x) \in \mathbb{R}^{2}: t>0 ; \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

with $\varphi(0)>0$ and consider the following problems:

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \text { a.e. on } D_{2}  \tag{2.13}\\
\left.u\right|_{\Gamma_{1}}=\left.u\right|_{\Gamma_{0}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \text { a.e. on } D_{2}  \tag{2.14}\\
\left.v\right|_{\Gamma_{2}}=\left.v\right|_{\Gamma_{0}}=0 \\
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=0
\end{array}\right.
$$

where $f_{i} \in L^{2}\left(D_{2}\right), i=1,2$ and $\Gamma_{0}$ is the part of $\partial D_{2}$ where $t=0$.


Fig.3: The unbounded domain $D_{2}$.

The change of variables indicated in the previous section transforms $D_{2}$ into the half strip $P=] 0,+\infty[\times] 0,1[$. So Problem (2.13) (respectively, Problem (2.14)) can be written as follows:

$$
\left\{\begin{array}{l}
\partial_{t} u+a(t, x) \partial_{x} u-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u=f_{1} \text { a.e. on } P,  \tag{2.15}\\
\left.u\right|_{t=0}=\left.u\right|_{x=0}=0 \\
\partial_{x} u+\left.\beta_{2} \varphi(t) u\right|_{x=1}=0,
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v+a(t, x) \partial_{x} v-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v=f_{2} \text { a.e. on } P,  \tag{2.16}\\
\left.v\right|_{t=0}=\left.v\right|_{x=1}=0 \\
\left.\partial_{x} v+\left.\beta_{1} \varphi(t) v\right|_{x=0}=0\right)
\end{array}\right.
$$

where $f_{i} \in L^{2}(P), i=1,2$ and the coefficients $a$ is that defined in Section 2.2.


Fig.4: The half strip $P$.

Let $f_{i}^{(n)}, i=1,2$ be the restriction $\left.f_{i}\right|_{]_{0, n[\times] 0,1}}, i=1,2$ for all $n \in \mathbb{N}^{*}$. Then, Theorem 2.2.1 shows that for all $n \in \mathbb{N}^{*}$, there exists a function $u_{n} \in \mathcal{H}^{1,2}\left(P_{n}\right)$ (respectively, $v_{n} \in \mathcal{H}^{1,2}\left(P_{n}\right)$ ) which solves the problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{n}+a(t, x) \partial_{x} u_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u_{n}=f_{1}^{(n)} \text { a.e. on } P_{n}  \tag{2.17}\\
\left.u_{n}\right|_{t=0}=\left.u_{n}\right|_{x=0}=0 \\
\partial_{x} u_{n}+\left.\beta_{2} \varphi(t) u_{n}\right|_{x=1}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v_{n}+a(t, x) \partial_{x} v_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v_{n}=f_{2}^{(n)} \text { a.e. on } P_{n}  \tag{2.18}\\
\left.v_{n}\right|_{t=0}=\left.v_{n}\right|_{x=1}=0 \\
\left.\partial_{x} v_{n}+\left.\beta_{1} \varphi(t) v_{n}\right|_{x=0}=0\right)
\end{array}\right.
$$

where $f_{i}^{(n)} \in L^{2}\left(P_{n}\right), i=1,2$ and $\left.P_{n}=\right] 0, n[\times] 0,1[$.


Fig.5: The truncated half strip $P_{n}$.

Now, let us prove an "energy" type estimate for the solutions $u_{n}$ (respectively, $v_{n}$ ) which will allow us to solve Problem (2.15) (respectively, Problem (2.16)) and then equivalently Problem (2.13) (respectively, Problem (2.14)).

Proposition 2.3.1. There exists a constant $K>0$ independent of $n$ such that

$$
\begin{aligned}
& \left\|u_{n}\right\|_{\mathcal{H}^{1,2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}(P)}^{2}, \\
& \left\|v_{n}\right\|_{\mathcal{H}^{1,2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}(P)}^{2}
\end{aligned}
$$

In order to prove Proposition 2.3.1, we need the following result:
Lemma 2.3.1. There exists a constant $K$ independent of $n$ such that

$$
\begin{aligned}
& \left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}(P)}^{2} \\
& \left\|v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}(P)}^{2}
\end{aligned}
$$

Proof. The Poincaré inequality gives $\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq K\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}$ and $\left\|v_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq K\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}$. Now, we estimate the inner products $\left\langle f_{1}^{(n)}, u_{n}\right\rangle$ and $\left\langle f_{2}^{(n)}, v_{n}\right\rangle$ in $L^{2}\left(P_{n}\right)$.

1) Estimation of $\left\langle f_{1}^{(n)}, u_{n}\right\rangle$ :

$$
\begin{aligned}
\left\langle f_{1}^{(n)}, u_{n}\right\rangle= & \int_{P_{n}} u_{n} \partial_{t} u_{n} d t d x+\int_{P_{n}} a(t, x) u_{n} \partial_{x} u_{n} d t d x-\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} u_{n} \partial_{x}^{2} u_{n} d t d x \\
= & \int_{\partial P_{n}}\left[\frac{1}{2}\left|u_{n}\right|^{2} \nu_{t}+a(t, x) \frac{1}{2}\left|u_{n}\right|^{2} \nu_{x}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x} u_{n} \cdot u_{n} \nu_{x}\right] d \sigma \\
& +\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} u_{n}\right)^{2} d t d x-\frac{1}{2} \int_{P_{n}} \partial_{x} a(t, x)\left|u_{n}\right|^{2} d t d x
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the unit outward normal vector at the boundary of $P_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $P_{n}$ where $t=0$, we have $u_{n}=0$ and consequently the corresponding boundary integral vanishes. On the part of the boundary where $t=n$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral is the following:

$$
\int_{0}^{1} \frac{1}{2}\left(u_{n}\right)^{2}(n, x) d x
$$

On the part of the boundary where $x=0$, we have $\nu_{x}=-1, \nu_{t}=0$ and $u_{n}(t, 0)=0$. Consequently, the corresponding integral vanishes. On the part of the boundary where $x=1$, we have $\nu_{x}=1, \nu_{t}=0$ and

$$
\partial_{x} u_{n}(t, 1)+\beta_{2} \varphi(t) u_{n}(t, 1)=0
$$

Consequently, the corresponding integral is

$$
\int_{0}^{n} \frac{\left(2 c(t) \beta_{2}-\varphi_{2}^{\prime}(t)\right)}{2 \varphi(t)}\left(u_{n}\right)^{2}(t, 1) d t
$$

Finally,

$$
\begin{aligned}
\left\langle f_{1}^{(n)}, u_{n}\right\rangle= & \int_{\varphi_{1}(n)}^{\varphi_{2}(n)} \frac{1}{2}\left(u_{n}\right)^{2}(n, x) d x+\int_{0}^{n} \frac{\left(2 c(t) \beta_{2}-\varphi_{2}^{\prime}(t)\right)}{2 \varphi(t)}\left(u_{n}\right)^{2}(t, 1) d t+\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} u_{n}\right)^{2} d t d x \\
& +\frac{1}{2} \int_{P_{n}} \frac{\varphi^{\prime}(t)}{\varphi(t)}\left|u_{n}\right|^{2} d t d x .
\end{aligned}
$$

Thanks to the condition (2.6) and since the function $\varphi$ increases, we obtain

$$
\left\langle f_{1}^{(n)}, u_{n}\right\rangle \geq \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} u_{n}\right)^{2} d t d x \geq C\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Hence, for all $\epsilon>0$,

$$
\begin{aligned}
\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} & \leq \frac{1}{C}\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq \frac{1}{C \epsilon}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\frac{\epsilon}{C}\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
\end{aligned}
$$

By using the Poincaré inequality, we obtain

$$
\left(1-\frac{\epsilon}{C}\right)\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq \frac{1}{C \epsilon}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Choosing $\epsilon$ small enough in the previous inequality, we prove the existence of a constant $K$ such that

$$
\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Since

$$
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq\left\|f_{1}\right\|_{L^{2}(P)}^{2}
$$

we obtain

$$
\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}(P)}^{2}
$$

1) Estimation of $\left\langle f_{2}^{(n)}, v_{n}\right\rangle$ : We have

$$
\begin{aligned}
\left\langle f_{2}^{(n)}, v_{n}\right\rangle= & \int_{P_{n}} v_{n} \partial_{t} v_{n} d t d x+\int_{P_{n}} a(t, x) v_{n} \partial_{x} v_{n} d t d x-\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} v_{n} \partial_{x}^{2} v_{n} d t d x \\
= & \int_{\partial P_{n}}\left[\frac{1}{2}\left|v_{n}\right|^{2} \nu_{t}+a(t, x) \frac{1}{2}\left|v_{n}\right|^{2} \nu_{x}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x} v_{n} \cdot v_{n} \nu_{x}\right] d \sigma \\
& \int_{P_{n}}+\frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} v_{n}\right)^{2} d t d x-\frac{1}{2} \int_{P_{n}} \partial_{x} a(t, x)\left|v_{n}\right|^{2} d t d x,
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the unit outward normal vector at the boundary of $P_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $P_{n}$ where $t=0$, we have $v_{n}=0$ and consequently the corresponding boundary integral vanishes. On the part of the boundary where $t=n$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral is the following:

$$
\int_{0}^{1} \frac{1}{2}\left(v_{n}\right)^{2}(n, x) d x
$$

On the part of the boundary where $x=0$, we have $\nu_{x}=-1, \nu_{t}=0$ and

$$
\partial_{x} v_{n}(t, 0)+\beta_{1} \varphi(t) v_{n}(t, 0)=0
$$

Consequently, the corresponding integral is

$$
\int_{0}^{n} \frac{\left(-2 c(t) \beta_{1}+\varphi_{1}^{\prime}(t)\right)}{2 \varphi(t)}\left(v_{n}\right)^{2}(t, 0) d t
$$

On the part of the boundary where $x=1$, we have $\nu_{x}=1, \nu_{t}=0$ and $\left.v_{n}(t, 1)\right)=0$. Consequently, the corresponding integral vanishes. Finally,

$$
\begin{aligned}
\left\langle f_{2}^{(n)}, v_{n}\right\rangle= & \int_{\varphi_{1}(n)}^{\varphi_{2}(n)} \frac{1}{2}\left(v_{n}\right)^{2}(n, x) d x+\int_{0}^{n} \frac{\left(-2 c(t) \beta_{1}+\varphi_{1}^{\prime}(t)\right)}{2 \varphi(t)}\left(v_{n}\right)^{2}(t, 0) d t+\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} v_{n}\right)^{2} d t d x \\
& +\frac{1}{2} \int_{P_{n}} \frac{\varphi^{\prime}(t)}{\varphi(t)}\left|v_{n}\right|^{2} d t d x
\end{aligned}
$$

Thanks to the condition (2.6) and since the function $\varphi$ increases, we obtain

$$
\left\langle f_{2}^{(n)}, v_{n}\right\rangle \geq \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left(\partial_{x} v_{n}\right)^{2} d t d x \geq C\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Hence, for all $\epsilon>0$,

$$
\begin{aligned}
\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} & \leq \frac{1}{C}\left\|v_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq \frac{1}{C \epsilon}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\frac{\epsilon}{C}\left\|v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
\end{aligned}
$$

By using the Poincaré inequality, we obtain

$$
\left(1-\frac{\epsilon}{C}\right)\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq \frac{1}{C \epsilon}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Choosing $\epsilon$ small enough in the previous inequality, we prove the existence of a constant $K$ such that

$$
\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Since

$$
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq\left\|f_{2}\right\|_{L^{2}(P)}^{2}
$$

we obtain

$$
\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}(P)}^{2}
$$

Remark 2.3.1. Similar computations show that the same result holds true when we substitute the condition that $\varphi$ increases in a neighborhood of $+\infty$ by the following:

$$
\left|\varphi^{\prime}(t)\right| \varphi(t) \leq M c(t)
$$

## Proof of Proposition 2.3.1

Let us denote the inner product in $L^{2}\left(P_{n}\right)$ by $\langle.,$.$\rangle , and set L:=\partial_{t}+a(t, x) \partial_{x}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2}$, then we have

$$
\begin{aligned}
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}= & \left\langle\partial_{t} u_{n}+a(t, x) \partial_{x} u_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u_{n}, \partial_{t} u_{n}+a(t, x) \partial_{x} u_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u_{n}\right\rangle \\
= & \left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2 \int_{P_{n}} a \partial_{t} u_{n} \partial_{x} u_{n} d t d x \\
& -2 \int_{P_{n}} a \frac{c(t)}{\varphi^{2}(t)} \partial_{x} u_{n} \partial_{x}^{2} u_{n} d t d x-2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x
\end{aligned}
$$

and

$$
\begin{gathered}
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}=\quad\left\langle\partial_{t} v_{n}+a(t, x) \partial_{x} v_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v_{n}, \partial_{t} u_{n}+a(t, x) \partial_{x} v_{n}-\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v_{n}\right\rangle \\
=\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a \partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2 \int_{P_{n}} a \partial_{t} v_{n} \partial_{x} v_{n} d t d x \\
-2 \int_{P_{n}} a \frac{c(t)}{\varphi^{2}(t)} \partial_{x} v_{n} \partial_{x}^{2} v_{n} d t d x-2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x
\end{gathered}
$$

Observe that the coefficients $a$ and $\frac{c(t)}{\varphi^{2}(t)}$ are bounded. So, thanks to Lemma 3.3.2, for all $\epsilon>0$ we obtain

$$
\begin{aligned}
& \left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}-2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x \\
& \leq\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}+2\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a \frac{c(t)}{\varphi^{2}(t)} \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+K_{1}\left(1+\frac{2}{\epsilon}\right)\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq K_{\epsilon}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}
\end{aligned}
$$

and

$$
\begin{aligned}
& \left.\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\| \frac{c(t)}{\varphi^{2}(t)} \partial_{x}^{2} v_{n}\right) \|_{L^{2}\left(P_{n}\right)}^{2}-2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x \\
& \leq\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a \partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a \partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}+2\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a \frac{c(t)}{\varphi^{2}(t)} \partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+K_{1}\left(1+\frac{2}{\epsilon}\right)\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq K_{\epsilon}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\epsilon\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2},
\end{aligned}
$$

where $K_{\epsilon}$ and $K_{1}$ are constants independent of $n$. Consequently

$$
\begin{equation*}
(1-\epsilon)\left(\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}\right) \leq 2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x+K_{\epsilon}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \tag{2.19}
\end{equation*}
$$

and

$$
\begin{equation*}
(1-\epsilon)\left(\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}\right) \leq 2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x+K_{\epsilon}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(P_{n}\right)}^{2} \tag{2.20}
\end{equation*}
$$

Estimation of $2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x$ : We have

$$
\partial_{t} u_{n} \partial_{x}^{2} u_{n}=\partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right)-\frac{1}{2} \partial_{t}\left(\partial_{x} u_{n}\right)^{2} .
$$

Then

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x & =2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right) d t d x-\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t}\left(\partial_{x} u_{n}\right)^{2} d t d x \\
& =\int_{\partial P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left[-\left(\partial_{x} u_{n}\right)^{2} \nu_{t}+2 \partial_{t} u_{n} \partial_{x} u_{n} \nu_{x}\right] d \sigma+\int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} u_{n}\right)^{2} d t d x
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the outward normal vector at the boundary of $P_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $P_{n}$ where $t=0$, we have $u_{n}=0$ and consequently $\partial_{x} u_{n}=0$. The corresponding boundary integral vanishes. On the part of the boundary where $t=n$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
-\int_{0}^{1} \frac{c(n)}{\varphi^{2}(n)}\left(\partial_{x} u_{n}\right)^{2}(n, x) d x
$$

is negative. On the part of the boundary where $x=0$, we have $\nu_{x}=-1, \nu_{t}=0$ and $u_{n}(t, 0)=0$. Consequently, the corresponding integral vanishes. On the part of the boundary where $x=1$, we have $\nu_{x}=1, \nu_{t}=0$ and

$$
\partial_{x} u_{n}(t, 1)+\beta_{2} \varphi(t) u_{n}(t, 1)=0
$$

Consequently, the corresponding integral is

$$
\left.\int_{0}^{n} \frac{-2 \beta_{2} c(t)}{\varphi(t)}\right) \partial_{t} u_{n}(t, 1) u_{n}(t, 1) d t=\frac{-\beta_{2} c(n)}{\varphi(n)} u_{n}^{2}(n, 1)+\int_{0}^{n} \beta_{2}\left(\frac{c(t)}{\varphi(t)}\right)^{\prime} u_{n}^{2}(t, 1) d t
$$

which is negative thanks to the condition (2.8) and to the fact that $\beta_{2}>0$. Finally,

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x= & -\int_{0}^{1} \frac{c(n)}{\varphi^{2}(n)}\left(\partial_{x} u_{n}\right)^{2}(n, x) d x-\frac{\beta_{2} c(n)}{\varphi(n)} u_{n}^{2}(n, 1)+\int_{0}^{n} \beta_{2}\left(\frac{c(t)}{\varphi(t)}\right)^{\prime} u_{n}^{2}(t, 1) d t \\
& +\int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} u_{n}\right)^{2} d t d x
\end{aligned}
$$

Note that the functions $\frac{c(t)}{\varphi^{2}(t)}$ and $\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}$ are bounded. So, by using Lemma 2.3.1, we deduce

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x & \leq \int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} u_{n}\right)^{2} d t d x \\
& \leq K_{2}\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \\
& \leq K_{3}\left\|f_{1}\right\|_{L^{2}(P)}^{2}
\end{aligned}
$$

where $K_{2}$ and $K_{3}$ are constants independent of $n$. Consequently, Choosing $\epsilon=\frac{1}{2}$ in the relationship (2.19), we obtain

$$
\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}(P)}^{2}
$$

Estimation of $2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x$ : We have

$$
\partial_{t} v_{n} \partial_{x}^{2} v_{n}=\partial_{x}\left(\partial_{t} v_{n} \partial_{x} v_{n}\right)-\frac{1}{2} \partial_{t}\left(\partial_{x} v_{n}\right)^{2} .
$$

Then

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x & =2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{x}\left(\partial_{t} v_{n} \partial_{x} v_{n}\right) d t d x-\int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t}\left(\partial_{x} v_{n}\right)^{2} d t d x \\
& =\int_{\partial P_{n}} \frac{c(t)}{\varphi^{2}(t)}\left[-\left(\partial_{x} v_{n}\right)^{2} \nu_{t}+2 \partial_{t} v_{n} \partial_{x} v_{n} \nu_{x}\right] d \sigma+\int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} v_{n}\right)^{2} d t d x
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the outward normal vector at the boundary of $P_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $P_{n}$ where $t=0$, we have $v_{n}=0$ and consequently $\partial_{x} v_{n}=0$. The corresponding boundary integral vanishes. On the part of the boundary where $t=n$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
-\int_{0}^{1} \frac{c(n)}{\varphi^{2}(n)}\left(\partial_{x} v_{n}\right)^{2}(n, x) d x
$$

is negative. On the part of the boundary where $x=0$, we have $\nu_{x}=-1, \nu_{t}=0$ and

$$
\partial_{x} v_{n}(t, 0)+\beta_{1} \varphi(t) v_{n}(t, 0)=0 .
$$

Consequently, the corresponding integral is

$$
\left.\int_{0}^{n} \frac{2 \beta_{1} c(t)}{\varphi(t)}\right) \partial_{t} v_{n}(t, 0) v_{n}(t, 0) d t=\frac{\beta_{1} c(n)}{\varphi(n)} v_{n}^{2}(n, 0)-\int_{0}^{n} \beta_{1}\left(\frac{c(t)}{\varphi(t)}\right)^{\prime} v_{n}^{2}(t, 0) d t
$$

which is negative thanks to the conditions (2.8) and the fact that $\beta_{2}>0$. On the part of the boundary where $x=1$, we have $\nu_{x}=1, \nu_{t}=0$ and $v_{n}(t, 1)=0$. Consequently, the corresponding integral vanishes. Finally,

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x= & -\int_{0}^{1} \frac{c(n)}{\varphi^{2}(n)}\left(\partial_{x} v_{n}\right)^{2}(n, x) d x+\frac{\beta_{1} c(n)}{\varphi(n)} v_{n}^{2}(n, 0)-\int_{0}^{n} \beta_{1}\left(\frac{c(t)}{\varphi(t)}\right)^{\prime} v_{n}^{2}(t, 0) d t \\
& +\int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} v_{n}\right)^{2} d t d x
\end{aligned}
$$

Note that the functions $\frac{c(t)}{\varphi^{2}(t)}$ and $\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}$ are bounded. So, by using Lemma 2.3.1, we deduce

$$
\begin{aligned}
2 \int_{P_{n}} \frac{c(t)}{\varphi^{2}(t)} \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x & \leq \int_{P_{n}}\left(\frac{c(t)}{\varphi^{2}(t)}\right)^{\prime}\left(\partial_{x} v_{n}\right)^{2} d t d x \\
& \leq K_{2}\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \\
& \leq K_{3}\left\|f_{2}\right\|_{L^{2}(P)}^{2}
\end{aligned}
$$

where $K_{2}$ and $K_{3}$ are constants independent of $n$. Consequently, Choosing $\epsilon=\frac{1}{2}$ in the relationship (2.20), we obtain

$$
\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}(P)}^{2}
$$

This ends the proof of Proposition 2.3.1.
Remark 2.3.2. We obtain the solution $u$ of Problem (2.13) (respectively, $v$ of Problem (2.14)) by letting $n$ go to infinity in the previous proposition. The uniqueness can be proved as in Theorem 2.2.1.

Finally, we have proved the following Theorem:
Theorem 2.3.1. Problem (2.13) (respectively, Problem (2.14)) admits a (unique) solution $u \in \mathcal{H}^{1,2}\left(D_{2}\right)$ (respectively, $v \in \mathcal{H}^{1,2}\left(D_{2}\right)$ ).

### 2.4 The case of a small in time bounded triangular domain

Let $T$ be a small enough positive real number. We set

$$
D_{3}:=\left\{(t, x) \in \mathbb{R}^{2}: 0<t<T ; \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

with $\varphi(0)=0$ and consider the following problems:

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \text { a.e. on } D_{3}  \tag{2.21}\\
\left.u\right|_{\Gamma_{1}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\Gamma_{2}}=0
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \text { a.e. on } D_{3}  \tag{2.22}\\
\left.v\right|_{\Gamma_{2}}=0 \\
\partial_{x} v+\left.\beta_{1} v\right|_{\Gamma_{1}}=0
\end{array}\right.
$$

where $f_{i} \in L^{2}\left(D_{3}\right), i=1,2$. Set

$$
Q_{n}=\left\{(t, x) \in D_{3}: \frac{1}{n}<t<T\right\}, n \in \mathbb{N}^{*} \text { and } \frac{1}{n}<T
$$

For each $n \in \mathbb{N}^{*}$ such that $\frac{1}{n}<T$, we set $f_{i}^{(n)}=\left.f_{i}\right|_{Q_{n}} \in L^{2}\left(Q_{n}\right), i=1,2$ and denote by $u_{n} \in \mathcal{H}^{1,2}\left(Q_{n}\right)$ (respectively, $\left.v_{n} \in \mathcal{H}^{1,2}\left(Q_{n}\right)\right)$ the solution of the following problem:

$$
\left\{\begin{array}{l}
\partial_{t} u_{n}-c(t) \partial_{x}^{2} u_{n}=f_{1} \text { a.e. on } Q_{n},  \tag{2.23}\\
\left.u_{n}\right|_{t=\frac{1}{n}}=\left.u_{n}\right|_{x=\varphi_{1}(t)}=0 \\
\partial_{x} u_{n}+\left.\beta_{2} u_{n}\right|_{x=\varphi_{2}(t)}=0 .
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v_{n}-c(t) \partial_{x}^{2} v_{n}=f_{2} \text { a.e. on } Q_{n}  \tag{2.24}\\
\left.v_{n}\right|_{t=\frac{1}{n}}=\left.v_{n}\right|_{x=\varphi_{2}(t)}=0 \\
\left.\partial_{x} v_{n}+\left.\beta_{1} v_{n}\right|_{x=\varphi_{1}(t)}=0\right)
\end{array}\right.
$$

Such a solution exists by Theorem 2.2.1.
Proposition 2.4.1. There exists a constant $K>0$ independent of $n$ such that

$$
\begin{aligned}
& \left\|u_{n}\right\|_{\mathcal{H}^{1,2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}^{2}, \\
& \left\|v_{n}\right\|_{\mathcal{H}^{1,2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)}^{2} .
\end{aligned}
$$

Remark 2.4.1. Let $\epsilon>0$ be a real which we will choose small enough. The hypothesis (3.3) implies the existence of a real number $T>0$ small enough such that

$$
\begin{equation*}
\left|\varphi_{i}^{\prime}(t) \varphi(t)\right| \leq \epsilon, \text { for all } t \in(0, T), i=1,2 . \tag{2.25}
\end{equation*}
$$

In order to prove Proposition 2.4.1, we need some preliminary results.
Lemma 2.4.1. There exists a constant $K$ independent of $n$ such that for all $t \in] 0, T[$ :

1) $\left\|u_{n}\right\|_{L^{2}\left(Q_{n}\right)} \leq K\left\|\varphi \partial_{x} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}, \quad\left\|v_{n}\right\|_{L^{2}\left(Q_{n}\right)} \leq K\left\|\varphi \partial_{x} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}$;
2) $\int_{\varphi_{1}(t)}^{\varphi_{2}(t)} u_{n}^{2}(t, x) d x \leq K \varphi^{4} \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left(\partial_{x}^{2} u_{n}\right)^{2}(t, x) d x, \quad \int_{\varphi_{1}(t)}^{\varphi_{2}(t)} v_{n}^{2}(t, x) d x \leq K \varphi^{4} \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left(\partial_{x}^{2} v_{n}\right)^{2}(t, x) d x$;
3) $\left.\int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left(\partial_{x} u_{n}\right)^{2}(t, x) d x \leq K \varphi^{2} \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left(\partial_{x}^{2} u_{n}\right)^{2}(t, x) d x, \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left(\partial_{x} v_{n}\right)^{2}(t, x) d x \leq K \varphi^{2} \int_{\varphi_{1}(t)}^{\varphi_{2}(t)} \partial_{x}^{2} v_{n}\right)^{2}(t, x) d x$;
4) $\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(Q_{n}\right)} \leq K\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}, \quad\left\|\partial_{x} v_{n}\right\|_{L^{2}\left(Q_{n}\right)} \leq K\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)}$.

Proof. Inequalities (1) are consequences of the Poincaré inequality.
The following operators are isomorphisms (see, [15])

$$
H_{\gamma}^{2}(0,1) \longrightarrow L^{2}(0,1), u \mapsto u^{\prime \prime}, \quad H_{\delta}^{2}(0,1) \longrightarrow L^{2}(0,1), v \mapsto v^{\prime \prime}
$$

where,

$$
H_{\gamma}^{2}(0,1)=\left\{u \in H^{2}(0,1): u(0)=0, u^{\prime}(1)+\beta_{2} u(1)=0\right\}
$$

and

$$
H_{\delta}^{2}(0,1)=\left\{v \in H^{2}(0,1): v^{\prime}(0)+\beta_{1} v(0)=0, v(1)=0\right\} .
$$

So, there exists a constant $K>0$ such that

$$
\begin{aligned}
&\|u\|_{L^{2}(0,1)} \leq\left\|u^{\prime \prime}\right\|_{L^{2}(0,1)},\|v\|_{L^{2}(0,1)} \\
&\left\|u^{\prime}\right\|_{L^{2}(0,1)} \leq\left\|v^{\prime \prime}\right\|_{L^{2}(0,1)} \\
& \leq\left\|u_{L^{2}(0,1)},\right\| v^{\prime} \|_{L^{2}(0,1)}
\end{aligned} \leq\left\|v^{\prime \prime}\right\|_{L^{2}(0,1)} .
$$

The change of variables (for a fixed $t$ )

$$
[0,1] \rightarrow\left[\varphi_{1}(t), \varphi_{2}(t)\right] ; x \longmapsto y=(1-x) \varphi_{1}(t)+x \varphi_{2}(t),
$$

leads to the estimates (2) and (3).
To prove (4), it is sufficient to expand the inner products $\left\langle f_{1}^{(n)}, u_{n}\right\rangle$ and $\left\langle f_{2}^{(n)}, v_{n}\right\rangle$ and use the inequalities (1). Indeed, we deduce for all $\epsilon>0$, (see the proof of uniqueness of solutions in Theorem 2.2.1)

$$
\begin{aligned}
\int_{Q_{n}} c(t)\left(\partial_{x} u_{n}\right)^{2} d t d x & \leq\left|\left\langle f_{1}^{(n)}, u_{n}\right\rangle\right| \\
& \leq \frac{1}{\epsilon}\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\epsilon\left\|u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \\
& \leq \frac{1}{\epsilon}\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}^{2}+\epsilon K\left\|\varphi u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}
\end{aligned}
$$

and

$$
\begin{aligned}
\int_{Q_{n}} c(t)\left(\partial_{x} v_{n}\right)^{2} d t d x & \leq\left|\left\langle f_{2}^{(n)}, v_{n}\right\rangle\right| \\
& \leq \frac{1}{\epsilon}\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\epsilon\left\|v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \\
& \leq \frac{1}{\epsilon}\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)}^{2}+\epsilon K\left\|\varphi v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}
\end{aligned}
$$

However, $\varphi$ is bounded and $c>\alpha>0$. Choosing $\epsilon$ small enough yields the desired result.
Proof of Proposition 2.4.1: Let us denote the inner product in $L^{2}\left(Q_{n}\right)$ by $\langle.,$.$\rangle and set$ $\mathcal{L}:=\partial_{t}-c(t) \partial_{x}^{2}$, then we have

$$
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2}=\left\langle\mathcal{L} u_{n}, \mathcal{L} u_{n}\right\rangle=\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|c(t) \partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}-2\left\langle\partial_{t} u_{n}, c(t) \partial_{x}^{2} u_{n}\right\rangle
$$

and

$$
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2}=\left\langle\mathcal{L} v_{n}, \mathcal{L} v_{n}\right\rangle=\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|c(t) \partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}-2\left\langle\partial_{t} v_{n}, c(t) \partial_{x}^{2} v_{n}\right\rangle
$$

Estimation of $-2\left\langle\partial_{t} u_{n}, c(t) \partial_{x}^{2} u_{n}\right\rangle$ : We have

$$
\partial_{t} u_{n} \partial_{x}^{2} u_{n}=\partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right)-\frac{1}{2} \partial_{t}\left(\partial_{x} u_{n}\right)^{2}
$$

Then,

$$
\begin{aligned}
-2\left\langle\partial_{t} u_{n}, c(t) \partial_{x}^{2} u_{n}\right\rangle & =-2 \int_{Q_{n}} c(t) \partial_{t} u_{n} \partial_{x}^{2} u_{n} d t d x \\
& =-2 \int_{Q_{n}} c(t) \partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right) d t d x+\int_{Q_{n}} c(t) \partial_{t}\left(\partial_{x} u_{n}\right)^{2} d t d x \\
& =\int_{\partial Q_{n}} c(t)\left[\left(\partial_{x} u_{n}\right)^{2} \nu_{t}-2 \partial_{t} u_{n} \partial_{x} u_{n} \nu_{x}\right] d \sigma-\int_{Q_{n}} c^{\prime}(t)\left(\partial_{x} u_{n}\right)^{2} d t d x
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the unit outward normal vector at the boundary of $Q_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $Q_{n}$ where $t=\frac{1}{n}$, we have $u_{n}=0$ and consequently $\partial_{x} u_{n}=0$. The corresponding boundary integral vanishes. On the part of the boundary where $t=T$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
\int_{\varphi_{1}(T)}^{\varphi_{2}(T)} c(T)\left(\partial_{x} u_{n}\right)^{2} d x
$$

is nonnegative. On the parts of the boundary where $x=\varphi_{i}(t), i=1,2$, we have

$$
\nu_{x}=\frac{(-1)^{i}}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}, \nu_{t}=\frac{(-1)^{i+1} \varphi_{i}^{\prime}(t)}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}, u_{n}\left(t, \varphi_{1}(t)\right)=\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)+\beta_{2} u_{n}\left(t, \varphi_{2}(t)\right)=0 .
$$

Consequently, the corresponding integral is

$$
\begin{aligned}
& -\int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left[\partial_{x} u_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} d t-2 \int_{\frac{1}{n}}^{T} c(t) \partial_{t} u_{n}\left(t, \varphi_{2}(t)\right) . \partial_{x} u_{n}\left(t, \varphi_{2}(t)\right) d t \\
& -\int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left[\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)\right]^{2} d t
\end{aligned}
$$

By putting $h(t):=u_{n}\left(t, \varphi_{2}(t)\right), t \in\left[\frac{1}{n}, T\right]$, we obtain

$$
\partial_{t} u_{n}\left(t, \varphi_{2}(t)\right) \partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)=h^{\prime}(t) \partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)-\varphi_{2}^{\prime}(t)\left(\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)\right)^{2}
$$

So, by using the boundary conditions, we get

$$
\begin{array}{rl}
-2 \int_{\frac{1}{n}}^{T} & c(t) \partial_{t} u_{n}\left(t, \varphi_{2}(t)\right) \partial_{x} u_{n}\left(t, \varphi_{2}(t)\right) d t \\
& =-2 \int_{\frac{1}{n}}^{T} c(t) h^{\prime}(t) \partial_{x} u_{n}\left(t, \varphi_{2}(t)\right) d t+2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left(\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)\right)^{2} d t \\
& =2 \beta_{2} \int_{\frac{1}{n}}^{T} c(t) h^{\prime}(t) h(t) d t+2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left(\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)^{2} d t\right. \\
& =\beta_{2} \int_{\frac{1}{n}}^{T} c(t)\left(h(t)^{2}\right)^{\prime} d t+2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left(\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)\right)^{2} d t \\
& =\beta_{2} c(T)(h(T))^{2}-\beta_{2} \int_{\frac{1}{n}}^{T} c^{\prime}(t) u_{n}^{2}\left(t, \varphi_{2}(t) d t+2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left(\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)^{2} d t\right.\right.
\end{array}
$$

Observe that, thanks to the condition (2.8) and the fact that $\beta_{2}>0, c(t)>0$, we have

$$
\beta_{2} c(T)(h(T))^{2}-\beta_{2} \int_{\frac{1}{n}}^{T} c^{\prime}(t) u_{n}^{2}\left(t, \varphi_{2}(t) d t \geq 0\right.
$$

So, by setting

$$
\begin{aligned}
& I_{n, 1}=-\int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left[\partial_{x} u_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} d t \\
& I_{n, 2}=\int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left[\partial_{x} u_{n}\left(t, \varphi_{2}(t)\right)\right]^{2} d t
\end{aligned}
$$

we have

$$
\begin{equation*}
-2\left\langle\partial_{t} u_{n}, c(t) \partial_{x}^{2} u_{n}\right\rangle \geq-\left|I_{n, 1}\right|-\left|I_{n, 2}\right| \tag{2.26}
\end{equation*}
$$

Estimation of $-2\left\langle\partial_{t} v_{n}, c(t) \partial_{x}^{2} v_{n}\right\rangle$ : We have

$$
\partial_{t} v_{n} \partial_{x}^{2} v_{n}=\partial_{x}\left(\partial_{t} v_{n} \partial_{x} v_{n}\right)-\frac{1}{2} \partial_{t}\left(\partial_{x} v_{n}\right)^{2} .
$$

Then,

$$
\begin{aligned}
-2\left\langle\partial_{t} u_{n}, c(t) \partial_{x}^{2} u_{n}\right\rangle & =-2 \int_{Q_{n}} c(t) \partial_{t} v_{n} \partial_{x}^{2} v_{n} d t d x \\
& =-2 \int_{Q_{n}} c(t) \partial_{x}\left(\partial_{t} v_{n} \partial_{x} v_{n}\right) d t d x+\int_{Q_{n}} c(t) \partial_{t}\left(\partial_{x} v_{n}\right)^{2} d t d x \\
& =\int_{\partial Q_{n}} c(t)\left[\left(\partial_{x} v_{n}\right)^{2} \nu_{t}-2 \partial_{t} v_{n} \partial_{x} v_{n} \nu_{x}\right] d \sigma-\int_{Q_{n}} c^{\prime}(t)\left(\partial_{x} v_{n}\right)^{2} d t d x
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ are the components of the unit outward normal vector at the boundary of $Q_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the part of the boundary of $Q_{n}$ where $t=\frac{1}{n}$, we have $v_{n}=0$ and consequently $\partial_{x} v_{n}=0$. The corresponding boundary integral vanishes. On the part of the boundary where $t=T$, we have $\nu_{x}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
\int_{\varphi_{1}(T)}^{\varphi_{2}(T)} c(T)\left(\partial_{x} v_{n}\right)^{2} d x
$$

is nonnegative. On the parts of the boundary where $x=\varphi_{i}(t), i=1,2$, we have

$$
\nu_{x}=\frac{(-1)^{i}}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}, \nu_{t}=\frac{(-1)^{i+1} \varphi_{i}^{\prime}(t)}{\sqrt{1+\left(\varphi_{i}^{\prime}\right)^{2}(t)}}, v_{n}\left(t, \varphi_{2}(t)\right)=\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)+\beta_{1} v_{n}\left(t, \varphi_{1}(t)\right)=0
$$

Consequently, the corresponding integral is

$$
\begin{aligned}
& \int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left[\partial_{x} v_{n}\left(t, \varphi_{2}(t)\right)\right]^{2} d t+2 \int_{\frac{1}{n}}^{T} c(t) \partial_{t} v_{n}\left(t, \varphi_{1}(t)\right) \cdot \partial_{x} v_{n}\left(t, \varphi_{1}(t)\right) d t \\
& \int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left[\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} d t .
\end{aligned}
$$

By putting $k(t):=v_{n}\left(t, \varphi_{1}(t)\right), t \in\left[\frac{1}{n}, T\right]$, we obtain

$$
\partial_{t} v_{n}\left(t, \varphi_{1}(t)\right) \partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)=k^{\prime}(t) \partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)-\varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right)^{2}
$$

So, by using the boundary conditions, we get

$$
\begin{aligned}
2 \int_{\frac{1}{n}}^{T} c(t) & \partial_{t} v_{n}\left(t, \varphi_{1}(t)\right) \partial_{x} v_{n}\left(t, \varphi_{1}(t)\right) d t \\
& =2 \int_{\frac{1}{n}}^{T} c(t) k^{\prime}(t) \partial_{x} v_{n}\left(t, \varphi_{1}(t)\right) d t-2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right)^{2} d t \\
& =-2 \beta_{1} \int_{\frac{1}{n}}^{T} c(t) k^{\prime}(t) k(t) d t-2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right)^{2} d t \\
& =-\beta_{1} \int_{\frac{1}{n}}^{T} c(t)\left(k(t)^{2}\right)^{\prime} d t-2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right)^{2} d t \\
& =-\beta_{1} c(T)(k(T))^{2}+\beta_{1} \int_{\frac{1}{n}}^{T} c^{\prime}(t) v_{n}^{2}\left(t, \varphi_{1}(t) d t-2 \int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right)^{2} d t\right.
\end{aligned}
$$

Observe that, thanks to the condition (2.8) and the fact that $\beta_{1}<0, c(t)>0$, we have

$$
-\beta_{1} c(T)(k(T))^{2}+\beta_{1} \int_{\frac{1}{n}}^{T} c^{\prime}(t) v_{n}^{2}\left(t, \varphi_{1}(t) d t \geq 0\right.
$$

So, by setting

$$
\begin{aligned}
J_{n, 1} & =\int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left[\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} d t \\
J_{n, 2} & =-\int_{\frac{1}{n}}^{T} c(t) \varphi_{2}^{\prime}(t)\left[\partial_{x} v_{n}\left(t, \varphi_{2}(t)\right)\right]^{2} d t
\end{aligned}
$$

we have

$$
\begin{equation*}
-2\left\langle\partial_{t} v_{n}, c(t) \partial_{x}^{2} v_{n}\right\rangle \geq-\left|J_{n, 1}\right|-\left|J_{n, 2}\right| \tag{2.27}
\end{equation*}
$$

Estimation of $I_{n, k}, J_{n, k}, k=1,2$ :
Lemma 2.4.2. There exists a constant $K>0$ independent of $n$ such that

$$
\begin{aligned}
\max \left(\left|I_{n, 1}\right|,\left|I_{n, 2}\right|\right) & \leq K \epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \\
\max \left(\left|J_{n, 1}\right|,\left|J_{n, 2}\right|\right) & \leq K \epsilon\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}
\end{aligned}
$$

Proof. We convert the boundary integral $J_{n, 1}$ into a surface integral by setting

$$
\begin{aligned}
{\left[\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} } & =-\left.\frac{\varphi_{2}(t)-x}{\varphi_{2}(t)-\varphi_{1}(t)}\left[\partial_{x} v_{n}(t, x)\right]^{2}\right|_{x=\varphi_{1}(t)} ^{x=\varphi_{2}(t)} \\
& =-\int_{\varphi_{1}(t)}^{\varphi_{2}(t)} \frac{\partial}{\partial x}\left\{\frac{\varphi_{2}(t)-x}{\varphi(t)}\left[\partial_{x} v_{n}(t, x)\right]^{2}\right\} d x \\
& =-2 \int_{\varphi_{1}(t)}^{\varphi_{2}(t)} \frac{\varphi_{2}(t)-x}{\varphi(t)} \partial_{x} v_{n}(t, x) \partial_{x}^{2} v_{n}(t, x) d x+\int_{\varphi_{1}(t)}^{\varphi_{2}(t)} \frac{1}{\varphi(t)}\left[\partial_{x} v_{n}(t, x)\right]^{2} d x
\end{aligned}
$$

Then, we have

$$
\begin{aligned}
J_{n, 1} & =\int_{\frac{1}{n}}^{T} c(t) \varphi_{1}^{\prime}(t)\left[\partial_{x} v_{n}\left(t, \varphi_{1}(t)\right)\right]^{2} d t \\
& =\int_{Q_{n}} \frac{c\left(t \varphi_{1}^{\prime}(t)\right.}{\varphi(t)}\left(\partial_{x} v_{n}\right)^{2} d t d x-2 \int_{Q_{n}} \frac{\varphi_{2}(t)-x}{\varphi(t)} c(t) \varphi_{1}^{\prime}(t)\left(\partial_{x} v_{n}\right)\left(\partial_{x}^{2} v_{n}\right) d t d x
\end{aligned}
$$

Thanks to Lemma 2.4.1, we can write

$$
\int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left[\partial_{x} v_{n}(t, x)\right]^{2} d x \leq C[\varphi(t)]^{2} \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left[\partial_{x}^{2} v_{n}(t, x)\right]^{2} d x
$$

Therefore,

$$
\int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left[\partial_{x} v_{n}(t, x)\right]^{2} \frac{\left|\varphi_{1}^{\prime}\right|}{\varphi} d x \leq C\left|\varphi_{1}^{\prime}\right|[\varphi] \int_{\varphi_{1}(t)}^{\varphi_{2}(t)}\left[\partial_{x}^{2} v_{n}(t, x)\right]^{2} d x
$$

consequently,

$$
\left|J_{n, 1}\right| \leq C \int_{Q_{n}} c(t)\left|\varphi_{1}^{\prime}\right|[\varphi]\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x+2 \int_{Q_{n}} c(t)\left|\varphi_{1}^{\prime}\right|\left|\partial_{x} v_{n}\right|\left|\partial_{x}^{2} v_{n}\right| d t d x
$$

since $\left|\frac{\varphi_{2}(t)-x}{\varphi(t)}\right| \leq 1$. So, for all $\epsilon>0$, we have

$$
\left|J_{n, 1}\right| \leq C \int_{Q_{n}}\left|c(t) \varphi_{1}^{\prime}\right|[\varphi]\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x+\epsilon \int_{Q_{n}} c(t)\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x+\frac{1}{\epsilon} \int_{Q_{n}} c(t)\left(\varphi_{1}^{\prime}\right)^{2}\left(\partial_{x} v_{n}\right)^{2} d t d x
$$

Lemma 2.4.1 yields

$$
\frac{1}{\epsilon} \int_{Q_{n}} c(t)\left(\varphi_{1}^{\prime}\right)^{2}\left(\partial_{x} v_{n}\right)^{2} d t d x \leq C \frac{1}{\epsilon} \int_{Q_{n}} c(t)\left(\varphi_{1}^{\prime}\right)^{2}[\varphi]^{2}\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x
$$

Thus, there exists a constant $M>0$ independent of $n$ such that

$$
\begin{aligned}
\left|J_{n, 1}\right| & \leq C \int_{Q_{n}} c(t)\left[\left|\varphi_{1}^{\prime}\right||\varphi|+\frac{1}{\epsilon}\left(\varphi_{1}^{\prime}\right)^{2}|\varphi|^{2}\right]\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x+\epsilon \int_{Q_{n}} c(t)\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x \\
& \leq M \epsilon \int_{Q_{n}}\left(\partial_{x}^{2} v_{n}\right)^{2} d t d x
\end{aligned}
$$

because $\left|\varphi_{1}^{\prime} \varphi\right| \leq \epsilon$. The inequalities

$$
\left|I_{n, 1}\right| \leq K \epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2},\left|I_{n, 2}\right| \leq K \epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2},\left|J_{n, 2}\right| \leq K \epsilon\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2},
$$

can be proved by a similar argument.
Now, we can complete the proof of Proposition 2.4.1. Summing up the estimates (2.26), (2.27) and those of Lemma 2.4.2, we then obtain

$$
\begin{aligned}
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} & \geq\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|c(t) \partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}-K_{1} \epsilon\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \\
& \geq\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left(\alpha^{2}-K_{1} \epsilon\right)\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}
\end{aligned}
$$

and

$$
\begin{aligned}
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} & \geq\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|c(t) \partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}-K_{2} \epsilon\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \\
& \geq\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left(\alpha^{2}-K_{2} \epsilon\right)\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2},
\end{aligned}
$$

where $K_{i}, i=1,2$ are positive numbers. Then, it is sufficient to choose $\epsilon$ such that

$$
\min \left(\alpha^{2}-K_{1} \epsilon, \alpha^{2}-K_{2} \epsilon\right)>0
$$

to get a constant $K_{0}>0$ independent of $n$ such that

$$
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \geq K_{0}\left(\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}\right)
$$

and

$$
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \geq K_{0}\left(\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}\right) .
$$

But

$$
\begin{aligned}
\left\|f_{1}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)} & \leq\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}, \\
\left\|f_{2}^{(n)}\right\|_{L^{2}\left(Q_{n}\right)} & \leq\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)},
\end{aligned}
$$

then, there exists a constant $K>0$, independent of $n$ satisfying

$$
\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}^{2}
$$

and

$$
\left\|\partial_{t} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2}+\left\|\partial_{x}^{2} v_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)}^{2} .
$$

Consequently, making use of Lemma 2.4.1 and the previous estimates, then, there exists a constant $K>0$, independent of $n$ satisfying

$$
\left\|u_{n}\right\|_{\mathcal{H}^{1,2}\left(Q_{n}\right)}^{2} \leq C\left\|f_{1}\right\|_{L^{2}\left(D_{3}\right)}^{2}
$$

and

$$
\left\|v_{n}\right\|_{\mathcal{H}^{1,2}\left(Q_{n}\right)}^{2} \leq C\left\|f_{2}\right\|_{L^{2}\left(D_{3}\right)}^{2} .
$$

This ends the proof of Proposition 2.4.1. Finally, we have proved the following Theorem:
Theorem 2.4.1. Problem (2.21) (respectively, (2.22)) admits a (unique) solution $u \in \mathcal{H}^{1,2}\left(D_{3}\right)$ (respectively, $v \in \mathcal{H}^{1,2}\left(D_{3}\right)$ ).

Proof. We obtain the solution $u$ of Problem (2.21) (respectively, $v$ of Problem (2.22)) by letting $n$ go to infinity in the previous proposition. The uniqueness can be proved as in Theorem 2.2.1.

### 2.5 Back to Problems (2.1)-(2.3) and (2.2)-(2.4) and proof of Theorem 2.1.1

The proof of Theorem 2.1.1 can be obtained by subdividing the domain

$$
D:=\left\{(t, x) \in \mathbb{R}^{2}: t>0, \varphi_{1}(t)<x<\varphi_{2}(t)\right\}
$$

into three open subdomains $\Omega_{1}, \Omega_{2}$ and $\Omega_{3}$. So, we set $D=\Omega_{1} \cup \Omega_{2} \cup \Omega_{3} \cup \Gamma_{T_{1}} \cup \Gamma_{T_{2}}$ where

$$
\begin{gathered}
\Omega_{1}=\left\{(t, x) \in D: 0<t<T_{1}\right\}, \Omega_{2}=\left\{(t, x) \in D: T_{1}<t<T_{2}\right\}, \Omega_{3}=\left\{(t, x) \in D: t>T_{2}\right\} \\
\Gamma_{T_{1}}=\left\{\left(T_{1}, x\right) \in \mathbb{R}^{2}: \varphi_{1}\left(T_{1}\right)<x<\varphi_{2}\left(T_{1}\right)\right\} \text { and } \Gamma_{T_{2}}=\left\{\left(T_{2}, x\right) \in \mathbb{R}^{2}: \varphi_{1}\left(T_{2}\right)<x<\varphi_{2}\left(T_{2}\right)\right\}
\end{gathered}
$$

with $T_{1}$ is a small enough positive number and $T_{2}$ is an arbitrary positive number such that $T_{2}>T_{1}$. In the sequel, $f_{1}, f_{2}$ stands for an arbitrary fixed elements of $L^{2}(D)$ and $\left(f_{1}^{(i)}, f_{2}^{(i)}\right)=$ $\left(\left.f_{1}\right|_{\Omega_{i}},\left.f_{2}\right|_{\Omega_{i}}\right), i=1,2,3$.

Theorem 3.4.1 applied to the triangular domain $\Omega_{1}$, shows that there exists a unique solution $w_{1} \in \mathcal{H}^{1,2}\left(\Omega_{1}\right)$ (respectively, $\left.w_{2} \in \mathcal{H}^{1,2}\left(\Omega_{1}\right)\right)$ of the problem

$$
\left\{\begin{array}{l}
\partial_{t} w_{1}-c(t) \partial_{x}^{2} w_{1}=f_{1}^{(1)} \quad \text { a.e. on } \Omega_{1}  \tag{2.28}\\
\left.w_{1}\right|_{\Gamma_{1,1}}=0 \\
\partial_{x} w_{1}+\left.\beta_{2} w_{1}\right|_{\Gamma_{2,1}}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} w_{2}-c(t) \partial_{x}^{2} w_{2}=f_{2}^{(1)} \text { a.e. on } \Omega_{1}  \tag{2.29}\\
\left.w_{2}\right|_{\Gamma_{2,1}}=0 \\
\left.\partial_{x} w_{2}+\left.\beta_{1} w_{2}\right|_{\Gamma_{1,1}}=0\right)
\end{array}\right.
$$

where $f_{i}^{(1)} \in L^{2}\left(\Omega_{1}\right), i=1,2$ and $\Gamma_{i, 1}$ are the parts of the boundary of $\Omega_{1}$ where $x=\varphi_{i}(t), i=$ 1,2 .

Lemma 2.5.1. If $w \in \mathcal{H}^{1,2}(] 0, T[\times] 0,1[)$, then $\left.w\right|_{t=0} \in H^{1}\left(\gamma_{0}\right),\left.w\right|_{x=0} \in H^{\frac{3}{4}}\left(\gamma_{1}\right)$ and $\left.w\right|_{x=1} \in$ $H^{\frac{3}{4}}\left(\gamma_{2}\right)$, where $\left.\gamma_{0}=\{0\} \times\right] 0,1\left[, \gamma_{1}=\right] 0, T\left[\times\{0\}\right.$ and $\left.\gamma_{2}=\right] 0, T[\times\{1\}$.

It is a particular case of Theorem 2.1 ([21], Vol.2). The transformation

$$
(t, x) \longmapsto\left(t^{\prime}, x^{\prime}\right)=\left(t, \varphi(t) x+\varphi_{1}(t)\right)
$$

leads to the following lemma:
Lemma 2.5.2. If $w \in \mathcal{H}^{1,2}\left(\Omega_{2}\right)$, then $\left.w\right|_{\Gamma_{T_{1}}} \in H^{1}\left(\Gamma_{T_{1}}\right),\left.w\right|_{x=\varphi_{1}(t)} \in H^{\frac{3}{4}}\left(\Gamma_{1,2}\right)$ and $\left.w\right|_{x=\varphi_{2}(t)} \in$ $H^{\frac{3}{4}}\left(\Gamma_{2,2}\right)$, where $\Gamma_{i, 2}$ are the parts of the boundary of $\Omega_{2}$ where $x=\varphi_{i}(t), i=1,2$.

Hereafter, we denote the trace $\left.w_{1}\right|_{\Gamma_{T_{1}}}$ (respectively, $\left.w_{2}\right|_{\Gamma_{T_{1}}}$ ) by $\psi_{1}$ (respectively, $\psi_{2}$ ) which is in the Sobolev space $H^{1}\left(\Gamma_{T_{1}}\right)$ because $w_{1} \in \mathcal{H}^{1,2}\left(\Omega_{1}\right)$ (respectively, $w_{2} \in \mathcal{H}^{1,2}\left(\Omega_{1}\right)$ ) (see Lemma 2.5.2). Now, consider the following problem in $\Omega_{2}$ :

$$
\left\{\begin{array}{l}
\partial_{t} w_{3}-c(t) \partial_{x}^{2} w_{3}=f_{1}^{(2)} \text { a.e. on } \Omega_{2}  \tag{2.30}\\
\left.w_{3}\right|_{\Gamma_{T_{1}}}=\psi_{1} \\
\left.w_{3}\right|_{\Gamma_{1,2}}=0 \\
\partial_{x} w_{3}+\left.\beta_{2} w_{3}\right|_{\Gamma_{2,2}}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} w_{4}-c(t) \partial_{x}^{2} w_{4}=f_{2}^{(2)} \text { a.e. on } \Omega_{2}  \tag{2.31}\\
\left.w_{4}\right|_{\Gamma_{T_{1}}}=\psi_{2} \\
\left.w_{4}\right|_{\Gamma_{2,2}}=0 \\
\partial_{x} w_{4}+\left.\beta_{1} w_{4}\right|_{\Gamma_{1,2}}=0
\end{array}\right.
$$

where $f_{i}^{(2)} \in L^{2}\left(\Omega_{2}\right), i=1,2$ and $\Gamma_{i, 2}$ are the parts of the boundary of $\Omega_{2}$ where $x=\varphi_{i}(t)$, $i=1,2$. We use the following result, which is a consequence of Theorem 4.3 ([21], Vol.2), to solve Problem (2.30).

Proposition 2.5.1. Let $Q$ be the rectangle $] 0, T[\times] 0,1\left[, f_{1}, f_{2} \in L^{2}(Q)\right.$ and $\psi_{1}, \psi_{2} \in H^{1}\left(\gamma_{0}\right)$. Then, the following problem admits a (unique) solution $u \in \mathcal{H}^{1,2}(Q)$ (respectively, $v \in \mathcal{H}^{1,2}(Q)$ ):

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \partial_{x}^{2} u=f_{1} \in L^{2}(Q) \\
\left.u\right|_{\gamma_{0}}=\psi_{1} \\
\left.u\right|_{\gamma_{1}}=0 \\
\partial_{x} u+\left.\beta_{2} u\right|_{\gamma_{2}}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} v-c(t) \partial_{x}^{2} v=f_{2} \in L^{2}(Q) \\
\left.v\right|_{\gamma_{0}}=\psi_{2} \\
\left.v\right|_{\gamma_{2}}=0 \\
\left.\partial_{x} v+\left.\beta_{1} v\right|_{\gamma_{1}}=0\right)
\end{array}\right.
$$

where $\left.\gamma_{0}=\{0\} \times\right] 0,1\left[, \gamma_{1}=\right] 0, T\left[\times\{0\}\right.$ and $\left.\gamma_{2}=\right] 0, T[\times\{1\}$.

Thanks to the transformation

$$
(t, x) \longmapsto(t, y)=\left(t, \varphi(t) x+\varphi_{1}(t)\right),
$$

we deduce the following result:
Proposition 2.5.2. Problem (2.30) (respectively, (2.31)) admits a (unique) solution $w_{3} \in$ $\mathcal{H}^{1,2}\left(\Omega_{2}\right)$ (respectively, $w_{4} \in \mathcal{H}^{1,2}\left(\Omega_{2}\right)$ ).

Hereafter, we denote the trace $\left.w_{3}\right|_{\Gamma_{T_{2}}}$ by $\Phi_{1}$ (respectively, $\left.w_{4}\right|_{\Gamma_{T_{2}}}$ by $\Phi_{2}$ ) which is in the Sobolev space $H^{1}\left(\Gamma_{T_{2}}\right)$ because $w_{3} \in \mathcal{H}^{1,2}\left(\Omega_{2}\right)$ (respectively, $w_{4} \in \mathcal{H}^{1,2}\left(\Omega_{2}\right)$ ) (see Lemma 2.5.2). Now, consider the following problem in $\Omega_{3}$ :

$$
\left\{\begin{array}{l}
\partial_{t} w_{5}-c(t) \partial_{x}^{2} w_{5}=f_{1}^{(3)} \text { a.e. on } \Omega_{3}  \tag{2.32}\\
\left.w_{5}\right|_{\Gamma_{T_{2}}}=\Phi_{1} \\
\left.w_{5}\right|_{\Gamma_{1,3}}=0 \\
\partial_{x} w_{5}+\left.\beta_{2} w_{5}\right|_{\Gamma_{2,3}}=0
\end{array}\right.
$$

(respectively,

$$
\left\{\begin{array}{l}
\partial_{t} w_{6}-c(t) \partial_{x}^{2} w_{6}=f_{2}^{(3)} \text { a.e. on } \Omega_{3},  \tag{2.33}\\
\left.w_{6}\right|_{\Gamma_{T_{2}}}=\Phi_{2} \\
\left.w_{6}\right|_{\Gamma_{2,3}}=0, \\
\left.\partial_{x} w_{6}+\left.\beta_{1} w_{6}\right|_{\Gamma_{1,3}}=0,\right)
\end{array}\right.
$$

where $f_{i}^{(3)} \in L^{2}\left(\Omega_{3}\right), i=1,2$ and $\Gamma_{i, 3}$ are the parts of the boundary of $\Omega_{3}$ where $x=\varphi_{i}(t)$, $i=1,2$. By similar arguments like those used previously, we deduce the following result:

Proposition 2.5.3. Problem (2.32) (respectively, (2.33)) admits a (unique) solution $w_{5} \in$ $\mathcal{H}^{1,2}\left(\Omega_{3}\right)$ (respectively, $w_{6} \in \mathcal{H}^{1,2}\left(\Omega_{3}\right)$ ).

Finally, the function $u$ (respectively, $v$ ) defined by

$$
u:=\left\{\begin{array}{l}
w_{1} \text { in } \Omega_{1} \\
w_{3} \text { in } \Omega_{2} \\
w_{5} \text { in } \Omega_{3}
\end{array}\right.
$$

(respectively,

$$
v:=\left\{\begin{array}{l}
w_{2} \text { in } \Omega_{1} \\
w_{4} \text { in } \Omega_{2} \\
w_{6} \text { in } \Omega_{3}
\end{array}\right.
$$

is the (unique) solution of Problem (2.1)-(2.3) (respectively, (2.2)-(2.4)). This ends the proof of Theorem 2.1.1.

## CHAPTER 3

## RESOLUTION OF A PARABOLIC EQUATION IN UNBOUNDED CONICAL DOMAINS OF $\mathbb{R}^{3}$

### 3.1 Introduction and statement of the main result

Let $D$ be an open set of $\mathbb{R}^{3}$ defined by

$$
D=\left\{(t, x, y) \in \mathbb{R}^{3}: t>0 ; 0 \leq \sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

where $\varphi \in C\left(\left[0,+\infty[) \cap C^{1}(0,+\infty)\right.\right.$,

$$
\varphi(t)>0 \quad \forall t>0, \text { and } \varphi(0)=0
$$

Let us introduce the following functional space

$$
\mathcal{H}^{1,2}(D)=\left\{u \in H^{1}(D): \partial_{x}^{2} u \in L^{2}(D), \partial_{y}^{2} u \in L^{2}(D), \partial_{x y}^{2} u \in L^{2}(D)\right\}
$$

where $H^{1}(D)$ stands for the Sobolev space defined by

$$
H^{1}(D)=\left\{u \in L^{2}(D): \partial_{t} u \in L^{2}(D), \partial_{x} u \in L^{2}(D), \partial_{y} u \in L^{2}(D)\right\}
$$

with $L^{2}(D)$ stands for the usual Lebesgue space of square-integrable functions on $D$. The space $\mathcal{H}^{1,2}(D)$ is equipped with the natural norm, that is

$$
\|u\|_{\mathcal{H}^{1,2}(D)}=\left(\|u\|_{H^{1}(D)}^{2}+\left\|\partial_{x}^{2} u\right\|_{L^{2}(D)}^{2}+\left\|\partial_{x y}^{2} u\right\|_{L^{2}(D)}^{2}+\left\|\partial_{y}^{2} u\right\|_{L^{2}(D)}^{2}\right)^{1 / 2}
$$

We consider the problem: to find a function $u \in \mathcal{H}^{1,2}(D)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } D \tag{3.1}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial D}=0, \tag{3.2}
\end{equation*}
$$

where $f \in L^{2}(D)$ and the coefficient $c$ is a continuous real-valued function defined on $[0,+\infty[$, differentiable on $] 0,+\infty[$ and such that

$$
0<\alpha \leq c(t) \leq \beta
$$

for every $t \in[0,+\infty[$, where $\alpha$ and $\beta$ are positive constants.


Fig.1: The unbounded conical domain D.

Observe that the main difficulties related to this kind of problems are due to the facts that $\varphi(0)=0$ and the unboundedness of $D$ with respect to the time variable. It is well known that there are two main approaches for the study of our problem. We can look for boundary conditions assuring the existence of the solution in the natural space, or we can work directly in the 'bad' domain which generates some singularities in the solution (see, for example [25] and
[12]). It is the first approach that we follow in this work. So, we impose sufficient conditions on the function $\varphi$ in the neighborhoods of 0 and $+\infty$. Our main result is the following :

Theorem 3.1.1. Let us assume that

$$
\begin{gather*}
\left|\varphi^{\prime}(t)\right| \varphi(t) \rightarrow 0 \quad \text { as } t \rightarrow 0^{+}  \tag{3.3}\\
\varphi \text { and } \varphi^{\prime} \text { are bounded in a neighborhood of }+\infty, \tag{3.4}
\end{gather*}
$$

and one of the following conditions is satisfied
(a) $\varphi$ is increasing in a neighborhood of $+\infty$,
(b) $\exists M>0:\left|\varphi^{\prime}\right| \varphi \leq M c(t)$.

Then Problem (3.1),(3.2) admits a unique solution $u \in \mathcal{H}^{1,2}(D)$.

### 3.2 Problem (3.1),(3.2) in a bounded domain which can be transformed into a cylinder

Let $T$ be an arbitrary positive number. Denote by

$$
D_{1}:=\left\{(t, x, y) \in \mathbb{R}^{3}: 0<t<T ; 0<\sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

with $\varphi(t)>0$ for all $t \in[0, T]$ and consider the following problem: to find a function $u \in$ $\mathcal{H}^{1,2}\left(D_{1}\right)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } D_{1} \tag{3.5}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial D_{1} \backslash \Gamma_{T}}=0, \tag{3.6}
\end{equation*}
$$

where $f \in L^{2}\left(D_{1}\right)$ and $\Gamma_{T}$ is the part of $\partial D_{1}$ where $t=T$, defined by

$$
\Gamma_{T}:=\left\{(T, x, y) \in \mathbb{R}^{3}: 0<\sqrt{x^{2}+y^{2}}<\varphi(T)\right\} .
$$



Fig.2: The bounded domain $D_{1}$.

Let us denote the inner product in $L^{2}\left(D_{1}\right)$ by $\langle.,$.$\rangle . Then, the uniqueness of the solutions$ may be obtained by developing the inner product $\left\langle\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right), u\right\rangle$. Indeed, Let us consider $u \in \mathcal{H}^{1,2}\left(D_{1}\right)$ a solution of Problem (3.5),(3.6) with a null right-hand side term. So,

$$
\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=0 \text { in } D_{1} .
$$

In addition $u$ fulfils the boundary condition (3.6). Using Green's formula, we have

$$
\begin{aligned}
& \int_{D_{1}}\left[\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)\right] u d t d x d y \\
& =\int_{\partial D_{1}}\left[\frac{1}{2}|u|^{2} \nu_{t}-c(t)\left(u . \partial_{x} u \nu_{x}+u . \partial_{y} u \nu_{y}\right)\right] d \sigma \\
& +\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}+\left|\partial_{x} u\right|^{2}\right) d t d x d y
\end{aligned}
$$

where $\nu_{t}, \nu_{x}, \nu_{y}$ are the components of the unit outward normal vector at $\partial D_{1}$. Taking into account the boundary conditions, all the boundary integrals vanish except $\int_{\partial D_{1}}|u|^{2} \nu_{t} d \sigma$. We have

$$
\int_{\partial D_{1}}|u|^{2} \nu_{t} d \sigma=\int_{\Gamma_{T}}|u|^{2} d x d y
$$

Then

$$
\begin{aligned}
& \int_{D_{1}}\left[\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)\right] u d t d x d y \\
& =\int_{\Gamma_{T}} \frac{1}{2}|u|^{2} d x d y+\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}+\left|\partial_{y} u\right|^{2}\right) d t d x d y
\end{aligned}
$$

Consequently

$$
\int_{D_{1}}\left[\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)\right] u d t d x d y=0
$$

yields

$$
\int_{D_{1}} c(t)\left(\left|\partial_{x} u\right|^{2}+\left|\partial_{y} u\right|^{2}\right) d t d x d y=0
$$

because

$$
\int_{\Gamma_{T}} \frac{1}{2}|u|^{2} d x d y \geq 0
$$

This implies that $\left|\partial_{x} u\right|^{2}+\left|\partial_{y} u\right|^{2}=0$, since $c(t)>0$ for all $t \in[0, T]$. Consequently, $\partial_{x}^{2} u+\partial_{y}^{2} u=0$. Then, the hypothesis $\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=0$ gives $\partial_{t} u=0$. Thus, $u$ is constant. The boundary conditions imply that $u=0$ in $D_{1}$. This proves the uniqueness of the solution of Problem (3.5),(3.6).

Now, let us look at the existence of solutions for Problem (3.5),(3.6). The change of variables $(t, x, y)$ to $\left(t, \frac{x}{\varphi(t)}, \frac{y}{\varphi(t)}\right)$ transforms $D_{1}$ into the cylinder $\left.Q=\right] 0, T[\times D(0,1)$, where $D(0,1)$ is the unit disk of $\mathbb{R}^{2}$ and Problem (3.5),(3.6) becomes the following: to find a function $u \in \mathcal{H}^{1,2}(Q)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u+\left(a(t, x) \partial_{x} u+a(t, y) \partial_{y} u\right)-b(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } Q \tag{3.7}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial Q \backslash(\{T\} \times D(0,1))}=0, \tag{3.8}
\end{equation*}
$$

where $f \in L^{2}(Q)$ and

$$
\begin{aligned}
& a(t, x)=-\frac{\varphi^{\prime}(t)}{\varphi(t)} x, a(t, y)=-\frac{\varphi^{\prime}(t)}{\varphi(t)} y \\
& b(t)=\frac{c(t)}{\varphi^{2}(t)} .
\end{aligned}
$$

The coefficients $a(t, x)$ and $a(t, y)$ are bounded. So, the operator

$$
\left[a(t, x) \partial_{x}+a(t, y) \partial_{y}\right]: \mathcal{H}^{1,2}(Q) \longrightarrow L^{2}(Q)
$$

is compact. Hence, it is sufficient to study the following problem: to find a function $u \in \mathcal{H}^{1,2}(Q)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-b(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } Q \tag{3.9}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial Q \backslash(\{T\} \times D(0,1))}=0 . \tag{3.10}
\end{equation*}
$$

It is clear that Problem (3.9),(3.10) admits a (unique) solution $u \in \mathcal{H}^{1,2}(Q)$ because the coefficient $b$ satisfies the "uniform parabolicity" condition (see, for example, [1]). On other hand, it is easy to verify that the aforementioned change of variables conserves the spaces $L^{2}$ and $\mathcal{H}^{1,2}$. Consequently, we have proved the following theorem:

Theorem 3.2.1. Problem (3.5),(3.6) admits a (unique) solution $u \in \mathcal{H}^{1,2}\left(D_{1}\right)$.

### 3.3 Well-posedeness results for Problem (3.1),(3.2) in a domain which can be transformed into a half strip

In this case, we set

$$
D_{2}:=\left\{(t, x, y) \in \mathbb{R}^{3}: t>0 ; 0<\sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

with $\varphi(0)>0$ and consider the following problem: to find a function $u \in \mathcal{H}^{1,2}\left(D_{2}\right)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } D_{2} \tag{3.11}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial D_{2}}=0 \tag{3.12}
\end{equation*}
$$

where $f \in L^{2}\left(D_{2}\right)$.


Fig.3: The unbounded domain $D_{2}$.
3.3. Well-posedeness results for Problem (3.1),(3.2) in a domain which can be transformed into a half strip

The change of variables indicated in the previous section transforms $D_{2}$ into the half strip $P=] 0,+\infty\left[\times D(0,1)\right.$, where $D(0,1)$ is the unit disk of $\mathbb{R}^{2}$. So Problem (3.11)-(3.12) can be written as follows: to find a function $u \in \mathcal{H}^{1,2}(P)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u+\left(a(t, x) \partial_{x} u+a(t, y) \partial_{y} u\right)-b(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } P \tag{3.13}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial P}=0, \tag{3.14}
\end{equation*}
$$

where the coefficients $a(t, x), a(t, y)$ and $b(t)$ are those defined in Section 2. Let $f_{n}$ be the restriction $\left.f\right|_{]_{0, n[\times D(0,1)}}$ for all $n \in \mathbb{N}^{*}$. Then, Theorem 3.2 .1 shows that for all $n \in \mathbb{N}^{*}$, there exists a function $u_{n} \in \mathcal{H}^{1,2}\left(P_{n}\right)$ which solves the problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{n}+\left(a(t, x) \partial_{x} u_{n}+a(t, y) \partial_{y} u_{n}\right)-b(t)\left(\partial_{x}^{2} u_{n}+\partial_{y}^{2} u_{n}\right)=f_{n} \in L^{2}\left(P_{n}\right)  \tag{3.15}\\
\left.u_{n}\right|_{\partial P_{n} \backslash(\{n\} \times D(0,1))}=0
\end{array}\right.
$$

where $\left.P_{n}=\right] 0, n\left[\times D(0,1)\right.$. Now, let us prove an "energy" type estimate for the solutions $u_{n}$ which will allow us to solve Problem (3.15) and then equivalently Problem (3.11)-(3.12).

Proposition 3.3.1. There exists a constant $K$ independent of $n$ such that

$$
\left\|u_{n}\right\|_{\mathcal{H}^{1,2}\left(P_{n}\right)} \leq K\|f\|_{L^{2}(P)}
$$

In order to prove Proposition 3.3.1, we need the following lemmas:
Lemma 3.3.1. There exists a constant $C>0$ independent of $n$ such that

$$
\left\|\partial_{x y}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{x}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{y}^{2} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq C\left\|\Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

The previous lemma is a consequence of Lemma 1.2.1 and Grisvard-Looss [9, Theorem 2.2].
Lemma 3.3.2. There exists a constant $K$ independent of $n$ such that

$$
\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq K\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq K\|f\|_{L^{2}(P)}
$$

where $\nabla u_{n}=\left(\partial_{x} u_{n}, \partial_{y} u_{n}\right)$.

Proof. The Poincaré inequality gives $\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)}$. Moreover, by developing the inner product $\left\langle\partial_{t} u_{n}+\left(a(t, x) \partial_{x} u_{n}+a(t, y) \partial_{y} u_{n}\right)-b(t)\left(\partial_{x}^{2} u_{n}+\partial_{y}^{2} u_{n}\right), u_{n}\right\rangle$ in $L^{2}\left(P_{n}\right)$, we obtain

$$
\begin{aligned}
\left\langle f_{n}, u_{n}\right\rangle & =\int_{P_{n}} u_{n}\left[\partial_{t} u_{n}+\left(a(t, x) \partial_{x} u_{n}+a(t, y) \partial_{y} u_{n}\right)-b(t)\left(\partial_{x}^{2} u_{n}+\partial_{y}^{2} u_{n}\right)\right] d t d x d y \\
& =\frac{1}{2} \int_{\Gamma_{n}}\left|u_{n}\right|^{2} d x d y+\int_{P_{n}} \frac{\varphi^{\prime}(t)}{\varphi(t)} u_{n}^{2}(t, x, y) d t d x d y+\int_{P_{n}} b(t)\left|\nabla u_{n}\right|^{2} d t d x d y .
\end{aligned}
$$

where $\Gamma_{n}=\left\{(n, x, y) \in \mathbb{R}^{3}: 0<\sqrt{x^{2}+y^{2}}<1\right\}$. Since the function $\varphi$ increases, we obtain

$$
\left\langle f_{n}, u_{n}\right\rangle \geq \int_{P_{n}} b(t)\left|\nabla u_{n}\right|^{2} d t d x d y \geq C\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
$$

Hence, for all $\epsilon>0$,

$$
\begin{aligned}
\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} & \leq \frac{1}{C}\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|f_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& \leq \frac{1}{C \epsilon}\|f\|_{L^{2}(P)}^{2}+\frac{\epsilon}{C}\left\|u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}
\end{aligned}
$$

By using the Poincaré inequality, we obtain

$$
\left(1-\frac{\epsilon}{C}\right)\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq \frac{1}{C \epsilon}\|f\|_{L^{2}(P)}^{2}
$$

Choosing $\epsilon$ small enough in the previous inequality, we prove the existence of a constant $K$ such that

$$
\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)} \leq K\|f\|_{L^{2}(P)}
$$

Remark 3.3.1. Similar computations show that the same result holds true when we substitute the condition that $\varphi$ increases in a neighborhood of $+\infty$ by the following:

$$
\left|\varphi^{\prime}(t)\right| \varphi(t) \leq M c(t)
$$

## Proof of Proposition 3.3.1: Set

$$
L u_{n}=\partial_{t} u_{n}+\left(a(t, x) \partial_{x} u_{n}+a(t, y) \partial_{y} u_{n}\right)-b(t) \Delta u_{n}
$$

where $\Delta u_{n}=\partial_{x}^{2} u_{n}+\partial_{y}^{2} u_{n}$. We have

$$
\begin{aligned}
\left\|f_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}= & \left\langle L u_{n}, L u_{n}\right\rangle \\
= & \left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a(t, x) \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a(t, y) \partial_{y} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|b(t) \Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \\
& +2 \int_{P_{n}} a(t, x) \partial_{t} u_{n} \partial_{x} u_{n} d t d x d y+2 \int_{P_{n}} a(t, y) \partial_{t} u_{n} \partial_{y} u_{n} d t d x d y \\
& -2 \int_{P_{n}} a(t, x) b(t) \partial_{x} u_{n} \Delta u_{n} d t d x d y-2 \int_{P_{n}} a(t, y) b(t) \partial_{y} u_{n} \Delta u_{n} d t d x d y \\
& -2 \int_{P_{n}} b \partial_{t} u_{n} \Delta u_{n} d t d x d y .
\end{aligned}
$$

Observe that the coefficients $a(t, x), a(t, y)$ and $b(t)$ are bounded. So, thanks to Lemma 3.3.2, for all $\epsilon>0$ we obtain

$$
\begin{aligned}
& \left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|b(t) \Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}-2 \int_{P_{n}} b \partial_{t} u_{n} \Delta u_{n} d t d x d y \\
\leq & \|f\|_{L^{2}(P)}^{2}+\left\|a(t, x) \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|a(t, y) \partial_{y} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a(t, x) \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& +2\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a(t, y) \partial_{y} u_{n}\right\|_{L^{2}\left(P_{n}\right)}+2\left\|\Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a(t, x) b(t) \partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
& +2\left\|\Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}\left\|a(t, y) b(t) \partial_{y} u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
\leq & \|f\|_{L^{2}(P)}^{2}+K_{1}\left(1+\frac{2}{\epsilon}\right)\left[\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\partial_{y} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}\right]+2 \epsilon\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2 \epsilon\left\|b(t) \Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)} \\
\leq & K_{\epsilon}\|f\|_{L^{2}(P)}^{2}+2 \epsilon\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+2 \epsilon\left\|\Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2},
\end{aligned}
$$

where $K_{\epsilon}$ and $K_{1}$ are constants independent of $n$. Consequently

$$
\begin{equation*}
(1-2 \epsilon)\left(\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|b(t) \Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}\right) \leq 2 \int_{P_{n}} b(t) \partial_{t} u_{n} \Delta u_{n} d t d x d y+K_{\epsilon}\|f\|_{L^{2}(P)}^{2} \tag{3.16}
\end{equation*}
$$

Estimation of $2 \int_{P_{n}} b(t) \partial_{t} u_{n} \Delta u_{n} d t d x d y$ : We have

$$
\partial_{t} u_{n} \cdot \Delta u_{n}=\partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right)+\partial_{y}\left(\partial_{t} u_{n} \partial_{y} u_{n}\right)-\frac{1}{2}\left[\partial_{t}\left(\partial_{x} u_{n}\right)^{2}+\partial_{t}\left(\partial_{y} u_{n}\right)^{2}\right] .
$$

Then

$$
\begin{aligned}
2 \int_{P_{n}} b(t) \partial_{t} u_{n} \Delta u_{n} d t d x d y= & 2 \int_{P_{n}} b(t) \partial_{x}\left(\partial_{t} u_{n} \partial_{x} u_{n}\right) d t d x d y+2 \int_{P_{n}} b(t) \partial_{y}\left(\partial_{t} u_{n} \partial_{y} u_{n}\right) d t d x d y \\
& -\int_{P_{n}} b(t) \partial_{t}\left(\partial_{x} u_{n}\right)^{2} d t d x d y-\int_{P_{n}} b(t) \partial_{t}\left(\partial_{y} u_{n}\right)^{2} d t d x d y \\
= & \int_{\partial P_{n}} b(t)\left[-\left|\nabla u_{n}\right|^{2} \nu_{t}+2 \partial_{t} u_{n}\left(\partial_{x} u_{n} \nu_{x}+\partial_{y} u_{n} \nu_{y}\right)\right] d \sigma \\
& +\int_{P_{n}} b^{\prime}(t)\left|\nabla u_{n}\right|^{2} d t d x d y
\end{aligned}
$$

where $\nu_{t}, \nu_{x}$ and $\nu_{y}$ are the components of the unit outward normal vector at $\partial P_{n}$. We shall rewrite the boundary integral making use of the boundary conditions. On the parts of the boundary of $P_{n}$ where $t=0$ and $\sqrt{x^{2}+y^{2}}=1$, we have $u_{n}=0$ and consequently $\partial_{x} u_{n}=$ $\partial_{y} u_{n}=0$. The corresponding boundary integral vanishes. On $\Gamma_{n}$, the part of the boundary where $t=n$, we have $\nu_{x}=\nu_{y}=0$ and $\nu_{t}=1$. Accordingly the corresponding boundary integral

$$
-\int_{\Gamma_{n}} b(n)\left|\nabla u_{n}\right|^{2}(n, x, y) d x d y
$$

is negative. Finally,

$$
2 \int_{P_{n}} b(t) \partial_{t} u_{n} \Delta u_{n} d t d x d y=-\int_{\Gamma_{n}} b(n)\left|\nabla u_{n}\right|^{2}(n, x, y) d x d y+\int_{P_{n}} b^{\prime}(t)\left|\nabla u_{n}\right|^{2} d t d x d y .
$$

Note that the functions $b(t)$ and $b^{\prime}(t)$ are bounded. So, by using Lemma 3.3.2, we deduce

$$
\begin{aligned}
2 \int_{P_{n}} b(t) \partial_{t} u_{n} \Delta u_{n} d t d x d y & \leq \int_{P_{n}} b^{\prime}(t)\left|\nabla u_{n}\right|^{2} d t d x d y \\
& \leq K_{2}\left\|\nabla u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \\
& \leq K_{3}\|f\|_{L^{2}(P)}^{2}
\end{aligned}
$$

where $K_{2}$ and $K_{3}$ are constants independent of $n$. Consequently, Choosing $\epsilon=\frac{1}{4}$ in the relationship (3.16), we obtain

$$
\left\|\partial_{t} u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2}+\left\|\Delta u_{n}\right\|_{L^{2}\left(P_{n}\right)}^{2} \leq K\|f\|_{L^{2}(P)}^{2} .
$$

This ends the proof of Proposition 3.3.1.
Remark 3.3.2. We obtain the solution $u$ of Problem (3.11)-(3.12) by letting $n$ go to infinity in the previous proposition. The uniqueness can be proved as in Theorem 3.2.1.

Finally, we have proved the following Theorem:
Theorem 3.3.1. Problem (3.11)-(3.12) admits a (unique) solution $u \in \mathcal{H}^{1,2}\left(D_{2}\right)$.

### 3.4 Local in time result

Let $T$ be a small enough positive real number. We set

$$
D_{3}:=\left\{(t, x, y) \in \mathbb{R}^{3}: 0<t<T ; 0 \leq \sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

with $\varphi(0)=0$ and consider the following problem: to find a function $u \in \mathcal{H}^{1,2}\left(D_{3}\right)$ that satisfies the equation

$$
\begin{equation*}
\partial_{t} u-c(t)\left(\partial_{x}^{2} u+\partial_{y}^{2} u\right)=f \text { a.e. on } D_{3} \tag{3.17}
\end{equation*}
$$

and the boundary condition

$$
\begin{equation*}
\left.u\right|_{\partial D_{3} \backslash \Gamma_{T}}=0, \tag{3.18}
\end{equation*}
$$

where $f \in L^{2}\left(D_{3}\right)$ and $\Gamma_{T}$ is the part of $\partial D_{3}$ where $t=T$, defined by

$$
\Gamma_{T}:=\left\{(T, x, y) \in \mathbb{R}^{3}: 0<\sqrt{x^{2}+y^{2}}<\varphi(T)\right\} .
$$

Set

$$
Q_{n}=\left\{(t, x, y) \in D_{3}: \frac{1}{n}<t<T\right\}, n \in \mathbb{N}^{*} \text { and } \frac{1}{n}<T
$$

For each $n \in \mathbb{N}^{*}$ such that $\frac{1}{n}<T$, we set $f_{n}=\left.f\right|_{Q_{n}} \in L^{2}\left(Q_{n}\right)$, and denote by $u_{n} \in \mathcal{H}^{1,2}\left(Q_{n}\right)$ the solution of the following problem:

$$
\left\{\begin{array}{l}
\partial_{t} u_{n}-c(t)\left(\partial_{x}^{2} u_{n}+\partial_{y}^{2} u_{n}\right)=f_{n} \text { a.e. on } Q_{n}  \tag{3.19}\\
\left.u_{n}\right|_{\partial Q_{n} \backslash \Gamma_{T}}=0
\end{array}\right.
$$

Such a solution exists by Theorem 3.2.1.
Proposition 3.4.1. There exists a constant $K>0$ independent of $n$ such that

$$
\left\|u_{n}\right\|_{\mathcal{H}^{1,2}\left(Q_{n}\right)}^{2} \leq K\left\|f_{n}\right\|_{L^{2}\left(Q_{n}\right)}^{2} \leq K\|f\|_{L^{2}\left(D_{3}\right)}^{2} .
$$

Remark 3.4.1. Let $\epsilon>0$ be a real which we will choose small enough. The hypothesis (3.3) implies the existence of a real number $T>0$ small enough such that

$$
\begin{equation*}
\left|\varphi^{\prime}(t) \varphi(t)\right| \leq \epsilon, \text { for all } t \in(0, T) \tag{3.20}
\end{equation*}
$$

In order to prove Proposition 3.4.1, we need some preliminary results.
Lemma 3.4.1. For a fixed $t \in] \frac{1}{n}, T\left[\right.$, let $\Omega_{t}$ be the bounded domain of $\mathbb{R}^{2}$ defined by

$$
\Omega_{t}=\left\{(x, y) \in \mathbb{R}^{2}: 0 \leq \sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

Then, there exists a constant $C>0$ such that
(a) $\max \left(\left\|\partial_{x} u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2},\left\|\partial_{y} u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2}\right) \leq C \varphi^{2}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2}$,
(b) $\left\|u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} \leq C \varphi^{4}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2}$.

Proof. It is a direct consequence of Lemma 1.2.1. Indeed, let $t \in] \frac{1}{n}, T$ [ and define the following change of variables

$$
\begin{aligned}
D(0,1) & \rightarrow \Omega_{t} \\
(x, y) & \longmapsto(\varphi(t) x, \varphi(t) y)=\left(x^{\prime}, y^{\prime}\right)
\end{aligned}
$$

Set

$$
v(x, y)=u_{n}(\varphi(t) x, \varphi(t) y)
$$

(a) We have

$$
\begin{aligned}
\left\|\partial_{x} v\right\|_{L^{2}(D(0,1))}^{2} & =\int_{D(0,1)}\left(\partial_{x} v\right)^{2}(x, y) d x d y \\
& =\int_{\Omega_{t}}\left(\partial_{x^{\prime}} u_{n}\right)^{2}\left(x^{\prime}, y^{\prime}\right) \varphi^{2}(t) \frac{1}{\varphi^{2}(t)} d x^{\prime} d y^{\prime} \\
& =\int_{\Omega_{t}}\left(\partial_{x^{\prime}} u_{n}\right)^{2}\left(x^{\prime}, y^{\prime}\right) d x^{\prime} d y^{\prime} \\
& =\left\|\partial_{x^{\prime}} u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
\end{aligned}
$$

On the other hand, we have

$$
\begin{aligned}
\|\Delta v\|_{L^{2}(D(0,1))}^{2} & =\int_{D(0,1)}\left[\left(\partial_{x}^{2} v+\partial_{y}^{2} v\right)(x, y)\right]^{2} d x d y \\
& =\int_{\Omega_{t}}\left[\varphi^{2}(t)\left(\partial_{x^{\prime}}^{2} u_{n}+\partial_{y^{\prime}}^{2} u_{n}\right)\right]^{2}\left(x^{\prime}, y^{\prime}\right) \frac{1}{\varphi^{2}(t)} d x^{\prime} d y^{\prime} \\
& =\varphi^{2}(t)\left[\partial_{x^{\prime}}^{2} u_{n}+\partial_{y^{\prime}}^{2} u_{n}\right]^{2}\left(x^{\prime}, y^{\prime}\right) d x^{\prime} d y^{\prime} \\
& =\varphi^{2}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
\end{aligned}
$$

Using the inequality

$$
\left\|\partial_{x} v\right\|_{L^{2}(D(0,1))}^{2} \leq C\|\Delta v\|_{L^{2}(D(0,1))}^{2}
$$

of Lemma 1.2.1, we obtain the desired inequality

$$
\left\|\partial_{x^{\prime}} u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} \leq C \varphi^{2}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
$$

By a similar argument, we get

$$
\left\|\partial_{y^{\prime}} u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} \leq C \varphi^{2}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
$$

(b) We have

$$
\|v\|_{L^{2}(D(0,1))}^{2}=\varphi^{-2}(t)\left\|u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
$$

On the other hand,

$$
\|\Delta v\|_{L^{2}(D(0,1))}^{2}=\varphi^{2}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
$$

Using the inequality

$$
\|v\|_{L^{2}(D(0,1))}^{2} \leq C\|\Delta v\|_{L^{2}(D(0,1))}^{2}
$$

of Lemma 1.2.1, we obtain the desired inequality

$$
\left\|u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} \leq C \varphi^{4}(t)\left\|\Delta u_{n}\right\|_{L^{2}\left(\Omega_{t}\right)}^{2} .
$$

Proof of Proposition 3.4.1: This result can be obtained by following step by step the proof of [[13], Proposition 3.2].

Theorem 3.4.1. Problem (3.17)-(3.18) admits a (unique) solution $u \in \mathcal{H}^{1,2}\left(D_{3}\right)$.
Proof. We obtain the solution $u$ of Problem (3.17)-(3.18) by letting $n$ go to infinity in the previous proposition.

### 3.5 Main result

The proof of Theorem 3.1.1 can be obtained by subdividing the domain

$$
D=\left\{(t, x, y) \in \mathbb{R}^{3}: t>0 ; 0 \leq \sqrt{x^{2}+y^{2}}<\varphi(t)\right\}
$$

into three subdomains $D_{1}, D_{2}$ and $D_{3}$. We set $D=D_{1} \cup D_{2} \cup D_{3} \cup \Gamma_{T_{1}} \cup \Gamma_{T_{2}}$, where

$$
\begin{gathered}
D_{1}=\left\{(t, x, y) \in D: 0<t<T_{1}\right\}, D_{2}=\left\{(t, x, y) \in D: T_{1}<t<T_{2}\right\} \\
D_{3}=\left\{(t, x, y) \in D: t>T_{2}\right\}, \Gamma_{T_{1}}:=\left\{\left(T_{1}, x, y\right) \in \mathbb{R}^{3}: 0<\sqrt{x^{2}+y^{2}}<\varphi\left(T_{1}\right)\right\},
\end{gathered}
$$

and

$$
\Gamma_{T_{2}}:=\left\{\left(T_{2}, x, y\right) \in \mathbb{R}^{3}: 0<\sqrt{x^{2}+y^{2}}<\varphi\left(T_{2}\right)\right\}
$$

with $T_{1}$ is a small enough positive number and $T_{2}$ is an arbitrary positive number such that $T_{2}>T_{1}$. In the sequel $f$ stands for an arbitrary fixed element of $L^{2}(D)$ and $f^{(i)}=f_{\left.\right|_{D_{i}}}, i=1,2,3$.

In the conical domain $D_{1}$, there exists a unique solution $u_{1} \in \mathcal{H}^{1,2}\left(D_{1}\right)$ of the problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{1}-c(t)\left[\partial_{x}^{2} u_{1}+\partial_{y}^{2} u_{1}\right]=f^{(1)} \in L^{2}\left(D_{1}\right)  \tag{3.21}\\
\left.u_{1}\right|_{\partial D_{1}-\Gamma_{T_{1}}}=0
\end{array}\right.
$$

Hereafter, we denote the trace $\left.u_{1}\right|_{\Gamma_{T_{1}}}$ by $\psi_{1}$ which is in the Sobolev space $H^{1}\left(\Gamma_{T_{1}}\right)$ because $u_{1} \in H^{1,2}\left(D_{1}\right)$ (see [21]). Now, consider the following problem in $D_{2}$

$$
\left\{\begin{array}{l}
\partial_{t} u_{2}-c(t)\left[\partial_{x}^{2} u_{2}+\partial_{y}^{2} u_{2}\right]=f^{(2)} \in L^{2}\left(D_{2}\right)  \tag{3.22}\\
\left.u_{2}\right|_{\Gamma_{T_{1}}}=\psi_{1}, \\
\left.u_{2}\right|_{\partial D_{2}-\left(\Gamma_{T_{1}} \cup \Gamma_{T_{2}}\right)}=0
\end{array}\right.
$$

We use the following result, which is a consequence of Theorem 4.3 ([21], Vol.2) to solve Problem (3.22).

Proposition 3.5.1. Let $Q$ be the cylinder $] 0, T\left[\times D(0,1)\right.$ where $D(0,1)$ is the unit disk of $\mathbb{R}^{2}$, $f \in L^{2}(Q)$ and $\phi \in H^{1}\left(\gamma_{0}\right)$. Then, the problem

$$
\left\{\begin{array}{l}
\partial_{t} u-c(t) \Delta u=f \text { in } Q, \\
\left.u\right|_{\gamma_{0}}=\phi, \\
\left.u\right|_{\gamma_{0} \cup \gamma_{1}}=0,
\end{array}\right.
$$

where $\left.\gamma_{0}=\{0\} \times D(0,1), \gamma_{1}=\right] 0, T\left[\times \partial D(0,1)\right.$, admits a (unique) solution $u \in H^{1,2}(Q)$.

Thanks to the transformation

$$
(t, x, y) \longmapsto\left(t, x^{\prime}, y^{\prime}\right)=(t, \varphi(t) x, \varphi(t) y),
$$

we deduce the following result :
Proposition 3.5.2. Problem (3.22) admits a (unique) solution $u_{2} \in \mathcal{H}^{1,2}\left(D_{2}\right)$.

Hereafter, we denote the trace $\left.u_{2}\right|_{\Gamma_{T_{2}}}$ by $\psi_{2}$ which is in the Sobolev space $H^{1}\left(\Gamma_{T_{2}}\right)$ because $u_{2} \in \mathcal{H}^{1,2}\left(D_{2}\right)$. Now, consider the following problem

$$
\left\{\begin{array}{l}
\partial_{t} u_{3}-c(t)\left[\partial_{x}^{2} u_{3}+\partial_{y}^{2} u_{3}\right]=f^{(3)} \in L^{2}\left(D_{3}\right)  \tag{3.23}\\
\left.u_{3}\right|_{\Gamma_{T_{2}}}=\psi_{2} \\
\left.u_{3}\right|_{\partial D_{3}-\Gamma_{T_{2}}}=0
\end{array}\right.
$$

By similar arguments like those used previously, we deduce the following result :
Proposition 3.5.3. Problem (3.23) admits a (unique) solution $u_{3} \in \mathcal{H}^{1,2}\left(D_{3}\right)$.

Finally, the function $u$ defined by 1

$$
u:=\left\{\begin{array}{lll}
u_{1} & \text { in } & D_{1} \\
u_{2} & \text { in } & D_{2} \\
u_{3} & \text { in } & D_{3}
\end{array}\right.
$$

is the (unique) solution of our problem. This ends the proof of Theorem 3.1.1.

## CONCLUSION AND PROSPECTS

There are many natural systems which evolve on time-dependent spatial domains. In this thesis, we have studied two boundary value problems associated to the following parabolic equation

$$
\partial_{t} u-c(t) \Delta u=f
$$

where $\Delta$ is the Laplacian operator, $c(t)$ is a time dependent coefficient and $f \in L^{2}$. The equation is posed in unbounded non regular domains and it is associated to different boundary conditions.

Whereas a comprehensive theory of such kind of problems exists in the case of bounded regular domains, the literature in the case of unbounded non-smooth domains does not seem to be very rich. One of the central results of this theory is the so-called "schift theorem": "If the second member is an element of a certain given regularity (in a Sobolev space for example), then the solution admits this regularity +2 ." In this thesis, we have given sufficient conditions on the functions of parametrization of the domains, and on the coefficients in order to bring to life this result when the domains are unbounded and non-smooth and the second member is in the Lebesgue space of square integrable functions, $L^{2}$.

First, we have proved well-posedness and regularity results for the above mentioned parabolic equation, subject to Dirichlet-Robin type boundary conditions and posed in an unbounded in
time non-rectangular domain. In the second work, the two dimensional case with CauchyDirichlet boundary conditions is studied in the case of unbounded conical domains.

In forthcoming works, the results obtained in this thesis will be extended at least in the following directions:

1. The function $f$ on the right-hand side of the above mentioned equation, may be taken in $L^{p}$, where $\left.p \in\right] 1,+\infty\left[\right.$ or in Holder spaces. The difficulty with the space $L^{p}, p \neq 2$, is that this space is not a Hilbert space, and so there is no an inner product. So, the domain decomposition method used in this cannot be generalized in this sense.
2. Consider higher-order parabolic equations and other boundary conditions.

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## ABSTRACT

This thesis is designed to give some results of existence and uniqueness of solutions for someparabolic equations posed in unbounded in time non-cylindrical domains. We give sufficient conditions on the functions of the parametrization of the non regular domains and on the coefficients of the equations under which our problems admit unique solutions. We study the global regularity problem in a suitable parabolic Sobolev space. The method used to prove our main results is based on the technique of the decomposition of domains.
Key words. parabolic equations, heat equation, non-rectangular domains, conical domains, unbounded domains, Dirichlet-Robin conditions, anisotropic Sobolev spaces.

## RÉSUMÉ

Cette thèse a pour but de donner des résultats d'existence et d'unicité de solutions pour certaines équations paraboliques posées dans des domaines non cylindriques et non bornés en temps. Nous donnons des conditions suffisantes sur les fonctions de paramétrisation des domaines non réguliers et sur les coefficients des équations sous lesquelles nos problèmes admettent une solution unique. Nous étudions le problème de régularité globale dans un espace de Sobolev parabolique approprié. La méthode utilisée pour démontrer nos principaux résultats est basée sur la technique de la décomposition de domaines.
Mots clés. Equations paraboliques, équation de la chaleur, domaines non rectangulaires, domaines coniques, domaines non bornés, conditions de Dirichlet-Robin, espaces de Sobolev anisotropes.

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الملخص
    تهوف هذه الاطروحة الى تقديم نتائج وجود ووحدانية الحلول لبعض المعادلات الدكافئية المطروحة في الساحات غير
    الأسطو انية غير المحدو دة بالنسبة للزمن. نقام شروطًا كافية على التو ابع المعلمة للساحات غير الصقيلة و على معاملات
المعادلات التي تحت ظرو فها تقبل مسائلنا حلوَا وحدانية. ندرس مسالة الصقالة الكلية في فضاءات سوبوليف غير متناظرة
    المناسبة. الطريقة المستخدمة لإثبات نتائجنا الرئيسية مبنية علىى تنتية تجزئة الساحات
    الكلمات الرئيسية. المعادلات الدكافئية، معادلة الحرارة، الساحات غير المستطيلة، الساحات المخروطية، الساحات غير
        المحدودة، شروط ديريشليه-روبن، فضاءات سوبوليف غير المثاظرة
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